



# Master loop equations, free energy and correlations for the chain of matrices.

Bertrand Eynard

## ► To cite this version:

Bertrand Eynard. Master loop equations, free energy and correlations for the chain of matrices.. 2003.  
hal-00000572v2

**HAL Id: hal-00000572**

**<https://hal.science/hal-00000572v2>**

Preprint submitted on 30 Oct 2003

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Master loop equations, free energy and correlations for the chain of matrices<sup>1</sup>

*\* Service de Physique Théorique, CEA/Saclay  
Orme des  
Merisiers F-91191 Gif-sur-Yvette Cedex, FRANCE*

## Abstract

The loop equations for a chain of hermitian random matrices are computed explicitly, including the  $1/N^2$  corrections. To leading order, the master loop equation reduces to an algebraic equation, whose solution can be written in terms of geometric properties of the underlying algebraic curve. In particular we compute the free energy, the resolvents, the 2-loop functions and some mixed one loop functions. We also initiate the calculation of the  $1/N^2$  expansion.

## 1 Introduction

The multimatrix model, or chain of matrices is defined by the probability weight for  $\mathcal{N} + 1$  hermitean matrices of size  $N \times N$ :

$$d\mu(M_0, \dots, M_{\mathcal{N}}) = \frac{1}{Z} dM_0 \dots dM_{\mathcal{N}} e^{-\frac{N}{T} \text{tr} [\sum_{k=0}^{\mathcal{N}} V_k(M_k) - \sum_{k=1}^{\mathcal{N}} M_{k-1} M_k]} \quad (1)$$

$$Z := (N/2\pi)^{-\frac{1}{2}(\mathcal{N}+1)N^2} e^{-\frac{N^2}{T^2}F} := \int dM_0 \dots dM_{\mathcal{N}} e^{-\frac{N}{T} \text{tr} [\sum_{k=0}^{\mathcal{N}} V_k(M_k) - \sum_{k=1}^{\mathcal{N}} M_{k-1} M_k]} \quad (2)$$

where  $M_k$  ( $k = 0, \dots, \mathcal{N}$ ) are  $N \times N$  hermitean matrices, and  $dM_k$  is the product of Lebesgue measures of all real components of  $M_k$ .  $Z$  is the partition function,  $F$  is the free energy.  $T$  is the temperature, it can be chosen equal

---

<sup>1</sup>Work partially supported by the EC IHP Network: “Discrete geometries: from solid state physics to quantum gravity”, HPRN-CT-1999-000161.

to 1 in most of the paper, except when one is interested in derivatives with respect to  $T$ .

The potentials  $V_k$  are polynomials of degree  $d_k + 1$ :

$$V_k(x) = g_{k,0} + \sum_{j=1}^{d_k+1} \frac{g_{k,j}}{j} x^j \quad (3)$$

In order for  $Z$  to exist, we have to assume that all  $V_k$ 's are bounded from below on the real axis. However, that constraint can be relaxed by studying "normal matrices" instead of hermitean matrices<sup>2</sup>, or by considering  $Z$  as a formal power series in the coefficients of the potentials. For simplicity, we assume here that the  $V_k$ 's are real<sup>3</sup>.

Each  $M_k$  lies in the potential well  $V_k$  and is linearly coupled to its neighbors  $M_{k-1}$  and  $M_{k+1}$ .  $M_0$  and  $M_{\mathcal{N}}$  have only one neighbor, and thus the chain is open. So far, the problem of the closed chain has remained unsolved.

The multimatrix model is a generalization of the 2-matrix model. It was often considered in the context of 2-dimensional quantum gravity and string theory. Its critical points are known to represent the minimal conformal field theories, characterized by a pair of integers  $(p, q)$ . It is known that one can get a  $(p, q)$  critical point with a multimatrix model where  $\mathcal{N} = q - 2$ . The necessity of studying multimatrix models can be understood from the fact that the one matrix models contains only the critical models with  $q \leq 2$ .

Recent progress have been made in the understanding of the two-matrix model [3, 4, 9, 10, 18, 19], and it has been often noticed that the chain of matrices presents many similarities with the 2-matrix model [7].

The loop equations of the 2-matrix model have been known for a long time [12, 22] and have been concisely written in [10] including the  $O(1/N^2)$  terms. The loop equations for the chain of matrices have been written in an appendix of [13] and in a draft of [15] without the  $O(1/N^2)$  terms, and without any proof, and also with misprints. The present paper is mainly a translation from french to english of the draft [15], with proofs, and additional results. It thus fills a gap in the appendix of [13]. In particular, the  $1/N^2$  terms are included, opening the way to finding next to leading order corrections

---

<sup>2</sup>i.e. the eigenvalues are located on complex paths instead of the real axis. For a potential of degree  $d_k + 1$ , there are  $d_k$  homologically non equivalent complex paths on which  $V_k$  is bounded from below [2].

<sup>3</sup>Most of the results presented here remain valid for complex  $V_k$ 's. The assumption of real potential merely allows to have a more concise derivation of loop equations.

as in [10]. The leading order solution is also presented in an updated algebraic geometry language.

The paper is organized as follows:

1. introduction
2. definitions of the loop functions.
3. the main result and its derivation: the master loop equation. Readers less interested in technical details can skip the derivation.
4. to leading order, the master loop equation is an algebraic equation. We discuss its solution.
5. computation of the free energy in the large N limit, along the lines drawn by [3].
6. computation of other resolvents and correlators in the large N limit.
7. next to leading order, along the lines drawn by [10, 11].
8. examples (2 and 3-matrix models, gaussian case, one-cut case).

The most important results of this paper are the master loop equation eq.(28), the free energy eq.(111), the two-loop function eq.(162), the mixed correlators eq.(165), eq.(182) and eq.(172).

## 2 Definitions

### 2.1 Moments

We define the moments:

$$T_{n_0, n_1, \dots, n_{\mathcal{N}}} := \frac{1}{2N} \langle \text{tr } M_0^{n_0} M_1^{n_1} \dots M_{\mathcal{N}}^{n_{\mathcal{N}}} \rangle + \frac{1}{2N} \langle \text{tr } M_{\mathcal{N}}^{n_{\mathcal{N}}} \dots M_1^{n_1} M_0^{n_0} \rangle \quad (4)$$

where  $\langle \rangle$  means averaging with the probability measure of eq.(1). If all the potentials  $V_k$ 's are bounded from below, it is clear that the moments are well defined convergent integrals. If on the contrary the  $V_k$ 's are not all bounded from below, then the moments are only formally defined through their expansions in the coupling constants.

## 2.2 The complete one loop-function

We define the complete one loop-function as follows:

$$W(z_0, \dots, z_{\mathcal{N}}) := \sum_{n_0, \dots, n_{\mathcal{N}}=0}^{\infty} \frac{T_{n_0, n_1, \dots, n_{\mathcal{N}}}}{z_0^{n_0+1} z_1^{n_1+1} \dots z_{\mathcal{N}}^{n_{\mathcal{N}}+1}} \quad (5)$$

which can also be written

$$\begin{aligned} W(z_0, \dots, z_{\mathcal{N}}) &= \frac{1}{2N} \left\langle \text{tr} \frac{1}{z_0 - M_0} \frac{1}{z_1 - M_1} \dots \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \right\rangle \\ &\quad + \frac{1}{2N} \left\langle \text{tr} \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \dots \frac{1}{z_1 - M_1} \frac{1}{z_0 - M_0} \right\rangle \end{aligned} \quad (6)$$

Notice that it is only a formal definition (even when all the potentials are bounded from below), it makes sense only through its power expansion in the vicinity of all  $z_k \rightarrow \infty$ . That function is merely a convenient concise notation for dealing with all the moments at once.

## 2.3 Uncomplete one loop-functions

We also define “uncomplete” loop-functions. For any integer  $k$  between zero and  $\mathcal{N} + 1$ , and for any ordered subset of  $[0, \mathcal{N}]$   $I = \{i_1, \dots, i_k\}$  ( $0 \leq i_1 < i_2 < \dots < i_k \leq \mathcal{N}$ ), we define:

$$\begin{aligned} W_{i_1, \dots, i_k}(z_{i_1}, \dots, z_{i_k}) &:= \frac{1}{2N} \left\langle \text{tr} \frac{1}{z_{i_1} - M_{i_1}} \frac{1}{z_{i_2} - M_{i_2}} \dots \frac{1}{z_{i_k} - M_{i_k}} \right\rangle \\ &\quad + \frac{1}{2N} \left\langle \text{tr} \frac{1}{z_{i_k} - M_{i_k}} \dots \frac{1}{z_{i_2} - M_{i_2}} \frac{1}{z_{i_1} - M_{i_1}} \right\rangle \\ &= (-1)^{\mathcal{N}+1-k} \text{Res} \prod_{j \notin I} dz_j W(z_0, \dots, z_{\mathcal{N}}) \\ &= \text{Pol}_{z_j, j \notin I} W(z_0, \dots, z_{\mathcal{N}}) \prod_{j \notin I} z_j \end{aligned} \quad (7)$$

where Pol means the polynomial part in the vicinity of  $\infty$ .

## 2.4 The resolvents

As a special case of eq.(7), the resolvent of the  $k^{\text{th}}$  matrix of the chain is defined by:

$$W_k(z_k) := \frac{1}{N} \left\langle \text{tr} \frac{1}{z_k - M_k} \right\rangle \quad (8)$$

In particular, the resolvent of the first matrix of the chain will play a prominent role:

$$W_0(z_0) := \frac{1}{N} \left\langle \text{tr} \frac{1}{z_0 - M_0} \right\rangle \quad (9)$$

The master loop equation derived in section 3.1 is an equation for  $W_0(z_0)$  as a function of  $z_0$ , and in section 4, we determine the large  $N$  limit of  $W_0(z_0)$ . In section 6.1, we derive the large  $N$  limits of the other resolvents  $W_k$ .

Instead of  $W_0(z_0)$ , it appears more convenient to consider the following function:

$$Z_1(z_0) := V'_0(z_0) - TW_0(z_0) \quad (10)$$

and more generally we define:

$$\begin{cases} Z_{-1}(z_0) := TW_0(z_0) \\ Z_0(z_0) := z_0 \\ Z_{k+1}(z_0) := V'_k(Z_k(z_0)) - Z_{k-1}(z_0) \end{cases} \quad 1 \leq k \leq \mathcal{N} \quad (11)$$

For short, we write:

$$Z_k := Z_k(z_0) \quad (12)$$

Notice that all  $Z_k$ 's are polynomials in both  $z_0$  and  $W_0(z_0)$ , and therefore, any polynomial of the  $Z_k$ 's is a polynomial in  $z_0$  and  $W_0(z_0)$ .

## 2.5 Polynomial loop functions $P$ and $U$

We define the functions  $f_{k,l}(z_k, \dots, z_l)$ :

$$\begin{cases} f_{k,l} := 0 & \text{if } l < k - 1 \\ f_{k,k-1} := 1 \\ f_{k,k}(z_k) := V'_k(z_k) \\ f_{k,l+1}(z_k, \dots, z_{l+1}) := V'_{l+1}(z_{l+1})f_{k,l}(z_k, \dots, z_l) - z_l z_{l+1} f_{k,l-1}(z_k, \dots, z_{l-1}) \end{cases} \quad (13)$$

Then we define for  $k \leq l$ :

$$P_{k,l}(z_0, \dots, z_{\mathcal{N}}) := \text{Pol}_{z_k, \dots, z_l} f_{k,l}(z_k, \dots, z_l) W(z_0, \dots, z_{\mathcal{N}}) \quad (14)$$

where  $\text{Pol}$  means the polynomial part in the vicinity of  $\infty$ .  $P_{k,l}$  depends on  $\mathcal{N} + 1$  variables, and is polynomial in the variables  $z_k, \dots, z_l$ .

By definition of  $f_{k,l}$  we have:

$$P_{k,l}(z_0, \dots, z_N) = \text{Pol}_{z_l} V'_l(z_l) P_{k,l-1}(z_0, \dots, z_N) - \text{Pol}_{z_l, z_{l-1}} z_{l-1} z_l P_{k,l-2}(z_0, \dots, z_N) \quad (15)$$

In particular, we define:

$$P(z_0, \dots, z_N) := P_{0,N}(z_0, \dots, z_N) \quad (16)$$

which is polynomial in all the variables, and

$$U(z_0, \dots, z_N) := P_{1,N}(z_0, \dots, z_N) \quad (17)$$

which is polynomial in all the variables but  $z_0$ .

We also define:

$$E(z_0, \dots, z_N) := (V'_0(z_0) - z_1)(V'_N(z_N) - z_{N-1}) - TP(z_0, \dots, z_N) \quad (18)$$

which is also a polynomial in all variables. These polynomials will play an important role in the master loop equation.

## 2.6 2-loop functions

We define the loop insertion operator as in [1]:

$$\frac{\partial}{\partial V_l(z)} := -\frac{1}{z} \frac{\partial}{\partial g_{l,0}} - \sum_{k=1}^{\infty} \frac{k}{z^{k+1}} \frac{\partial}{\partial g_{l,k}} \quad (19)$$

where  $V_l(z) = g_{l,0} + \sum_{k \geq 1} \frac{g_{l,k}}{k} z^k$ , and the derivatives are taken at  $g_{l,k} = 0$  for  $k > d_l + 1$ . It is a formal definition, which makes sense only order by order in its large  $z$  expansion.

It is such that (for  $|z| > |z'|$ ):

$$\frac{\partial V_k(z')}{\partial V_l(z)} = -\delta_{k,l} \frac{1}{z - z'}, \quad \frac{\partial V'_k(z')}{\partial V_l(z)} = -\delta_{k,l} \frac{1}{(z - z')^2} \quad (20)$$

When applied to eq.(2), it produces a resolvent, i.e. the expectation of a trace:

$$\frac{\partial F}{\partial V_l(z)} = -TW_l(z) = -\frac{T}{N} \left\langle \text{tr} \frac{1}{z - M_l} \right\rangle \quad (21)$$

and more generally, the action of  $\frac{\partial}{\partial V_l(z)}$  on an expectation value inserts a new trace.

In particular, we define the following two loop functions (i.e. two-traces):

$$\begin{aligned}
W_{;l}(z_0, \dots, z_{\mathcal{N}}; z) &:= T \frac{\partial}{\partial V_l(z)} W(z_0, \dots, z_{\mathcal{N}}) \\
&= \frac{1}{2} \left\langle \text{tr} \frac{1}{z_0 - M_0} \frac{1}{z_1 - M_1} \dots \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \text{tr} \frac{1}{z - M_l} \right\rangle \\
&\quad + \frac{1}{2} \left\langle \text{tr} \frac{1}{z - M_l} \text{tr} \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \dots \frac{1}{z_1 - M_1} \frac{1}{z_0 - M_0} \right\rangle \\
&\quad - N^2 W_l(z) W(z_0, \dots, z_{\mathcal{N}}) \tag{22}
\end{aligned}$$

$$\begin{aligned}
W_{i_1, \dots, i_k; l}(z_{i_1}, \dots, z_{i_k}; z) &:= T \frac{\partial}{\partial V_l(z)} W_{i_1, \dots, i_k}(z_{i_1}, \dots, z_{i_k}) \\
&= (-1)^{\mathcal{N}+1-k} \text{Res} \prod_{j \notin I} dz_j W_{;l}(z_0, \dots, z_{\mathcal{N}}; z) \tag{23}
\end{aligned}$$

$$P_{k,l;j}(z_0, \dots, z_{\mathcal{N}}; z) := \text{Pol}_{z_k, \dots, z_l} f_{k,l}(z_k, \dots, z_l) W_{;j}(z_0, \dots, z_{\mathcal{N}}; z) \tag{24}$$

$$P_{;j}(z_0, \dots, z_{\mathcal{N}}; z) := P_{0,\mathcal{N};j}(z_0, \dots, z_{\mathcal{N}}; z) \tag{25}$$

$$U_{;j}(z_0, \dots, z_{\mathcal{N}}; z) := P_{1,\mathcal{N};j}(z_0, \dots, z_{\mathcal{N}}; z) \tag{26}$$

In particular, the following two-loop functions will play an important role:

$$\begin{aligned}
W_{k;l}(z'; z) &= T \frac{\partial W_k(z')}{\partial V_l(z)} = - \frac{\partial^2 F}{\partial V_k(z') \partial V_l(z)} = T \frac{\partial W_l(z)}{\partial V_k(z')} = W_{l;k}(z; z') \\
&= \left\langle \text{tr} \frac{1}{z - M_l} \text{tr} \frac{1}{z' - M_k} \right\rangle_c \tag{27}
\end{aligned}$$

where the subscript c means connected part  $\langle AB \rangle_c := \langle AB \rangle - \langle A \rangle \langle B \rangle$ .

### 3 Loop equations

#### 3.1 Main result: the master loop equation

We prove below that the master loop equation (named after [22]) can be written:

$$E(Z_0, Z_1, \dots, Z_{\mathcal{N}}) = \frac{T^2}{N^2} U_{;0}(Z_0, Z_1, \dots, Z_{\mathcal{N}}; Z_0)$$

(28)



where the functions  $Z_k(z_0)$  have been defined in eq.(11). Examples for  $\mathcal{N} = 1$  (the 2-matrix model) and  $\mathcal{N} = 2$  (the 3-matrix model) are explicated in section 8.

As an intermediate step in proving eq.(28), we need to prove the following formula for all  $1 \leq k \leq \mathcal{N}$ :

$$(z_k - Z_k)W_{0,k,\dots,\mathcal{N}}(z_0, z_k, \dots, z_{\mathcal{N}}) = W_{0,k+1,\dots,\mathcal{N}}(z_0, z_{k+1}, \dots, z_{\mathcal{N}}) \\ - P_{0,k-1}(z_0, Z_1, \dots, Z_{k-1}, z_k, \dots, z_{\mathcal{N}}) \\ - \frac{T}{N^2} P_{1,k-1;0}(z_0, Z_1, \dots, Z_{k-1}, z_k, \dots, z_{\mathcal{N}}; z_0)$$

(29)

### 3.2 Derivation of the master loop equation

- Proof of eq.(29) for  $k = 1$ :

The invariance of the matrix integral under the infinitesimal change of variable (see [14, 8, 10], and pay attention to the non-commutativity of matrices, and to the hermiticity of the change of variables):

$$\delta M_0 = \frac{T}{2} \frac{1}{z_1 - M_1} \frac{1}{z_2 - M_2} \cdots \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \frac{1}{z_0 - M_0} \\ + \frac{T}{2} \frac{1}{z_0 - M_0} \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \cdots \frac{1}{z_2 - M_2} \frac{1}{z_1 - M_1} \quad (30)$$

implies

$$TW_0(z_0)W(z_0, z_1, \dots, z_{\mathcal{N}}) + \frac{T}{N^2} W_{;0}(z_0, z_1, \dots, z_{\mathcal{N}}; z_0) \\ = V'_0(z_0)W(z_0, z_1, \dots, z_{\mathcal{N}}) - P_{0,0}(z_0, z_1, \dots, z_{\mathcal{N}}) \\ - z_1 W_{0,1,\dots,\mathcal{N}}(z_0, z_1, \dots, z_{\mathcal{N}}) + W_{0,2,\dots,\mathcal{N}}(z_0, z_2, \dots, z_{\mathcal{N}}) \quad (31)$$

The LHS comes from the Jacobian of the change of variable, and the RHS comes from the variation of the action. The first two terms of the RHS come from  $\delta \text{tr } V_0(M_0) = \text{tr } V'_0(M_0) \delta M_0$ , and use of eq.(5). The last two terms of the RHS come from  $\delta \text{tr } M_0 M_1 = \text{tr } M_1 \delta M_0$ , and use of eq.(6).

Using  $Z_1 = V'_0(z_0) - TW_0(z_0)$ , this can be rewritten:

$$(z_1 - Z_1)W_{0,1,\dots,\mathcal{N}}(z_0, z_1, \dots, z_{\mathcal{N}}) = W_{0,2,\dots,\mathcal{N}}(z_0, z_2, \dots, z_{\mathcal{N}})$$

$$\begin{aligned}
& -P_{0,0}(z_0, z_1, \dots, z_{\mathcal{N}}) \\
& -\frac{T}{N^2} W_{;0}(z_0, z_1, \dots, z_{\mathcal{N}}; z_0)
\end{aligned}
\tag{32}$$

Therefore we have derived eq.(29) in the case  $k = 1$ .

- Proof of eq.(29) for  $k = 2$ :

Notice that eq.(32) implies:

$$\begin{aligned}
-\text{Res } dz_1 W(z_0, z_1, z_2, \dots, z_{\mathcal{N}}) V_1'(z_1) &= V_1'(Z_1) W_{0,2,\dots,\mathcal{N}}(z_0, z_2, \dots, z_{\mathcal{N}}) \\
& -P_{0,1}(z_0, Z_1, z_2, \dots, z_{\mathcal{N}}) \\
& -W_{2,\dots,\mathcal{N}}(z_2, \dots, z_{\mathcal{N}}) \\
& -\frac{T}{N^2} P_{1,1;0}(z_0, Z_1, z_2, \dots, z_{\mathcal{N}}; z_0)
\end{aligned}
\tag{33}$$

The change of variable:

$$\delta M_1 = \frac{1}{2} \frac{1}{z_2 - M_2} \cdots \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \frac{1}{z_0 - M_0} + \frac{1}{2} \frac{1}{z_0 - M_0} \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \cdots \frac{1}{z_2 - M_2}
\tag{34}$$

gives (since  $\delta M_1$  is independant of  $M_1$  there is no jacobian, the LHS is zero, and the RHS is the variation of  $\text{tr } V_1'(M_1) - M_0 M_1 - M_1 M_2$ ):

$$\begin{aligned}
0 &= -\text{Res } dz_1 W(z_0, z_1, z_2, \dots, z_{\mathcal{N}}) V_1'(z_1) \\
& -z_0 W_{0,2,\dots,\mathcal{N}}(z_0, z_2, \dots, z_{\mathcal{N}}) + W_{2,\dots,\mathcal{N}}(z_2, \dots, z_{\mathcal{N}}) \\
& -z_2 W_{0,2,\dots,\mathcal{N}}(z_0, z_2, \dots, z_{\mathcal{N}}) + W_{0,3,\dots,\mathcal{N}}(z_0, z_3, \dots, z_{\mathcal{N}})
\end{aligned}
\tag{35}$$

i.e., using eq.(33) and  $V_1'(Z_1) = z_0 + Z_2$ :

$$\begin{aligned}
(z_2 - Z_2) W_{0,2,\dots,\mathcal{N}}(z_0, z_2, \dots, z_{\mathcal{N}}) &= W_{0,3,\dots,\mathcal{N}}(z_0, z_3, \dots, z_{\mathcal{N}}) \\
& -P_{0,1}(z_0, Z_1, z_2, \dots, z_{\mathcal{N}}) \\
& -\frac{T}{N^2} P_{1,1;0}(z_0, Z_1, z_2, \dots, z_{\mathcal{N}}; z_0)
\end{aligned}
\tag{36}$$

Therefore we have derived eq.(29) in the case  $k = 2$ .

- Proof of eq.(29) by recursion on  $k$ :

We have already proved eq.(29) for  $k = 1$  and  $k = 2$ . Now assume that eq.(29) holds for  $k - 1$  and  $k$ , we are going to prove it for  $k + 1$ . Note that eq.(29) for  $k - 1$  implies (multiply by  $z_{k-1}$  and take the residue at  $z_{k-1} \rightarrow \infty$  and  $z_k \rightarrow \infty$ ):

$$\begin{aligned}
& \text{Res } dz_{k-1} \text{ Res } dz_k z_{k-1} W_{0,k-1,k,\dots,\mathcal{N}}(z_0, z_{k-1}, z_k, z_{k+1}, \dots, z_{\mathcal{N}}) \\
= & Z_{k-1} W_{0,k+1,\dots,\mathcal{N}}(z_0, z_{k+1}, \dots, z_{\mathcal{N}}) \\
& - \text{Pol}_{z_{k-1}, z_k} z_{k-1} z_k P_{0,k-2}(z_0, \dots, z_{\mathcal{N}}) \Big|_{z_0=Z_0, \dots, z_k=Z_k} \\
& - \frac{T}{N^2} \text{Pol}_{z_{k-1}, z_k} z_k z_{k-1} P_{1,k-2;0}(z_0, \dots, z_{\mathcal{N}}; z_0) \Big|_{z_0=Z_0, \dots, z_k=Z_k}
\end{aligned} \tag{37}$$

and eq.(29) for  $k$  implies (multiply by  $V'_k(z_k)$  and take the residue at  $z_k \rightarrow \infty$ ):

$$\begin{aligned}
& - \text{Res } dz_k V'_k(z_k) W_{0,k,k+1,\dots,\mathcal{N}}(z_0, z_k, z_{k+1}, \dots, z_{\mathcal{N}}) \\
= & V'_k(Z_k) W_{0,k+1,\dots,\mathcal{N}}(z_0, z_{k+1}, \dots, z_{\mathcal{N}}) \\
& - \text{Pol}_{z_k} V'_k(z_k) P_{0,k-1}(z_0, \dots, z_{\mathcal{N}}) \Big|_{z_0=Z_0, \dots, z_k=Z_k} \\
& - \frac{T}{N^2} \text{Pol}_{z_k} V'_k(z_k) P_{1,k-1;0}(z_0, \dots, z_{\mathcal{N}}; z_0) \Big|_{z_0=Z_0, \dots, z_k=Z_k}
\end{aligned} \tag{38}$$

Then consider the change of variable:

$$\begin{aligned}
\delta M_k = & \frac{1}{2} \frac{1}{z_{k+1} - M_{k+1}} \cdots \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \frac{1}{z_0 - M_0} \\
& + \frac{1}{2} \frac{1}{z_0 - M_0} \frac{1}{z_{\mathcal{N}} - M_{\mathcal{N}}} \cdots \frac{1}{z_{k+1} - M_{k+1}}
\end{aligned} \tag{39}$$

it gives

$$\begin{aligned}
0 = & - \text{Res } dz_k V'_k(z_k) W_{0,k,k+1,\dots,\mathcal{N}}(z_0, z_k, z_{k+1}, \dots, z_{\mathcal{N}}) \\
& - \text{Res } dz_{k-1} \text{ Res } dz_k z_{k-1} W_{0,k-1,k,\dots,\mathcal{N}}(z_0, z_{k-1}, \dots, z_{\mathcal{N}}) \\
& - z_{k+1} W_{0,k+1,\dots,\mathcal{N}}(z_0, z_{k+1}, \dots, z_{\mathcal{N}}) + W_{0,k+2,\dots,\mathcal{N}}(z_0, z_{k+2}, \dots, z_{\mathcal{N}})
\end{aligned} \tag{40}$$

using eq.(37) and eq.(38), as well as  $V'_k(Z_k) = Z_{k-1} + Z_{k+1}$ , we get eq.(29) for  $k + 1$ .

- eq.(29) for  $k = \mathcal{N}$ :

In particular, the previous recurrence derivation shows that eq.(29) for  $k = \mathcal{N}$  reads:

$$(z_{\mathcal{N}} - Z_{\mathcal{N}})W_{0,\mathcal{N}}(z_0, z_{\mathcal{N}}) = \begin{aligned} & W_0(z_0) \\ & - P_{0,\mathcal{N}-1}(z_0, Z_1, \dots, Z_{\mathcal{N}-1}, z_{\mathcal{N}}) \\ & - \frac{T}{N^2} P_{1,\mathcal{N}-1;0}(z_0, Z_1, \dots, Z_{\mathcal{N}-1}, z_{\mathcal{N}}; z_0) \end{aligned} \quad (41)$$

- Proof of eq.(28):

In particular, from eq.(29) for  $k = \mathcal{N} - 1$  we derive:

$$\begin{aligned} & \text{Res } dz_{\mathcal{N}-1} \text{Res } dz_{\mathcal{N}} z_{\mathcal{N}-1} W_{0,z_{\mathcal{N}-1},z_{\mathcal{N}}}(z_0, z_{\mathcal{N}-1}, z_{\mathcal{N}}) \\ = & Z_{\mathcal{N}-1} W_0(z_0) \\ & - \text{Pol}_{z_{\mathcal{N}-1},z_{\mathcal{N}}} z_{\mathcal{N}-1} z_{\mathcal{N}} P_{0,\mathcal{N}-2}(z_0, \dots, z_{\mathcal{N}}) \big|_{z_0=Z_0, \dots, z_{\mathcal{N}}=Z_{\mathcal{N}}} \\ & - \frac{T}{N^2} \text{Pol}_{z_{\mathcal{N}-1},z_{\mathcal{N}}} z_{\mathcal{N}-1} z_{\mathcal{N}} P_{1,\mathcal{N}-1;0}(z_0, \dots, z_{\mathcal{N}}; z_0) \big|_{z_0=Z_0, \dots, z_{\mathcal{N}}=Z_{\mathcal{N}}} \end{aligned} \quad (42)$$

and from eq.(41), we derive:

$$\begin{aligned} & - \text{Res } dz_{\mathcal{N}} V'_{\mathcal{N}}(z_{\mathcal{N}}) W_{0,\mathcal{N}}(z_0, z_{\mathcal{N}}) \\ = & V'_{\mathcal{N}}(Z_{\mathcal{N}}) W_0(z_0) \\ & - \text{Pol}_{z_{\mathcal{N}}} V'_{\mathcal{N}}(z_{\mathcal{N}}) P_{0,\mathcal{N}-1}(z_0, \dots, z_{\mathcal{N}}) \big|_{z_0=Z_0, \dots, z_{\mathcal{N}}=Z_{\mathcal{N}}} \\ & - \frac{T}{N^2} \text{Pol}_{z_{\mathcal{N}}} V'_{\mathcal{N}}(z_{\mathcal{N}}) P_{1,\mathcal{N}-1;0}(z_0, \dots, z_{\mathcal{N}}; z_0) \big|_{z_0=Z_0, \dots, z_{\mathcal{N}}=Z_{\mathcal{N}}} \end{aligned} \quad (43)$$

Then, the change of variables

$$\delta M_{\mathcal{N}} = \frac{1}{z_0 - M_0} \quad (44)$$

gives

$$\begin{aligned} 0 = & - \text{Res } dz_{\mathcal{N}} V'_{\mathcal{N}}(z_{\mathcal{N}}) W_{0,\mathcal{N}}(z_0, z_{\mathcal{N}}) \\ & - \text{Res } dz_{\mathcal{N}-1} \text{Res } dz_{\mathcal{N}} z_{\mathcal{N}-1} W_{0,z_{\mathcal{N}-1},z_{\mathcal{N}}}(z_0, z_{\mathcal{N}-1}, z_{\mathcal{N}}) \end{aligned} \quad (45)$$

i.e., using eq.(42) and eq.(43):

$$0 = (V'_{\mathcal{N}}(Z_{\mathcal{N}}) - Z_{\mathcal{N}-1}) W_0(z_0)$$

$$\begin{aligned}
& -P_{0,\mathcal{N}}(Z_0, \dots, Z_{\mathcal{N}}) \\
& -\frac{T}{N^2}P_{1,\mathcal{N};0}(Z_0, \dots, Z_{\mathcal{N}}; z_0)
\end{aligned}
\tag{46}$$

Recalling that  $Z_1 = V'_0(z_0) - TW_0(z_0)$ , we get:

$$\begin{aligned}
& (V'_0(Z_0) - Z_1)(V'_{\mathcal{N}}(Z_{\mathcal{N}}) - Z_{\mathcal{N}-1}) - TP_{0,\mathcal{N}}(Z_0, Z_1, \dots, Z_{\mathcal{N}}) \\
& = \frac{T^2}{N^2}P_{1,\mathcal{N};0}(Z_0, Z_1, \dots, Z_{\mathcal{N}}; Z_0)
\end{aligned}
\tag{47}$$

which ends the proof of eq.(28).

## 4 Large $N$ leading order, algebraic equation

Throughout this section and sections 5 and 6, we abusively denote with the same name, the loop functions and their large  $N$  limits.

To leading order, the master loop equation reduces to an algebraic equation for  $W_0$  as a function of  $z_0$ :

$$E(z_0, Z_1, \dots, Z_{\mathcal{N}}) = 0 \tag{48}$$

where  $Z_1 = V'_0(z_0) - TW_0(z_0)$  and  $Z_{k+1} = V'_k(Z_k) - Z_{k-1}$ .

$E$  is a polynomial of given degrees in each variable, and with known leading coefficients. The problem is that most of the subleading coefficients of  $E(z_0, z_1, \dots, z_{\mathcal{N}})$  are not determined by the loop equations. They are determined by additional hypothesis, which will be explained in section 4.5. Prior to that, we need to study the geometry of that algebraic equation.

The complex curve  $W_0$  (equivalently  $Z_1$ ) as a function of  $z_0$  is a one-dimensional submanifold of  $\mathbf{C}^{N+1}$ , which is the intersection of  $\mathcal{N}$  dimension  $\mathcal{N}$  algebraic submanifolds. The curve  $W_0$  as a function of  $z_0$  is thus a Riemann surface  $\mathcal{E}$ .

Instead of viewing the  $Z_k$ 's as (multivalued) functions of  $z_0$ , it is more appropriate to view the  $Z_k$ 's as (monovalued) complex functions over  $\mathcal{E}$ :

$$p \rightarrow z_k(p) \tag{49}$$

such that for all  $p \in \mathcal{E}$ :

$$Z_k(z_0(p)) = z_k(p) \tag{50}$$

Since the function  $z_0(p)$  is not injective, i.e. the point  $p$  such that  $z_0(p) = z_0$  is not unique, the  $Z_k$ 's are multivalued functions of  $z_0$ . Similarly, one could consider any  $z_k$  as a multivalued function of any  $z_l$ .

Instead of dealing with multivalued functions, we slice  $\mathcal{E}$  into domains called sheets, such that in each domain the functions we are considering are bijections. The  $z_0$ -sheets are thus domains in which the function  $z_0(p)$  is injective. More generally, the  $z_k$ -sheets are domains in which the function  $z_k(p)$  is injective.

Let us study the  $z_0$ -sheets first. For that purpose, we define:

$$\begin{cases} r_{-1} &:= -1 \\ r_0 &:= 1 \\ r_k &:= d_0 d_1 \dots d_{k-1}, \quad k = 1, \dots, \mathcal{N} + 1 \\ s_{\mathcal{N}+1} &:= -1 \\ s_{\mathcal{N}} &:= 1 \\ s_k &:= d_{k+1} \dots d_{\mathcal{N}-1} d_{\mathcal{N}}, \quad k = -1, \dots, \mathcal{N} - 1 \end{cases} \quad (51)$$

## 4.1 $z_0$ -sheets

The equation

$$E(z_0, Z_1, \dots, Z_{\mathcal{N}}) = 0 \quad (52)$$

has degree  $1 + d_1 d_2 \dots d_{\mathcal{N}}$  in  $W_0$  and degree  $d_0 d_1 \dots d_{\mathcal{N}}$  in  $z_0$ . Therefore,  $W_0(z_0)$  (or equivalently  $Z_1$ ) is a multivalued function of  $z_0$  which takes  $1 + d_1 d_2 \dots d_{\mathcal{N}}$  values: in other words, there are  $1 + d_1 d_2 \dots d_{\mathcal{N}}$   $z_0$ -sheets. We identify these sheets by looking at the asymptotics of  $Z_{\mathcal{N}}$  when  $z_0 \rightarrow \infty$ .

- The physical sheet

From eq.(9), there must exist at least one solution of the algebraic equation such that:

$$W_0(z_0) \underset{z_0 \rightarrow \infty}{\sim} \frac{1}{z_0} + O(1/z_0^2) \quad (53)$$

which implies that:

$$z_k = O(z_0^{r_k}) \quad (54)$$

and in particular

$$z_{\mathcal{N}} = O(z_0^{d_0 d_1 \dots d_{\mathcal{N}-1}}) \quad (55)$$

The  $z_0$ -sheet in which these asymptotics hold is called the  $z_0$ -physical sheet.

- Other sheets

Notice that the equation for  $W_{\mathcal{N}}(z_{\mathcal{N}})$  as a function of  $z_{\mathcal{N}}$  is the same algebraic equation. In other words, there is a solution of the algebraic equation which is such that:

$$z_k = O(z_{\mathcal{N}}^{s_k}) , \quad s_k = d_{k+1} \dots d_{\mathcal{N}-1} d_{\mathcal{N}} \quad (56)$$

i.e.

$$z_{\mathcal{N}} = O(z_0^{1/d_1 \dots d_{\mathcal{N}-1} d_{\mathcal{N}}}) , \quad z_k = O(z_0^{1/d_1 \dots d_k}) \quad (57)$$

Since the number of  $s_0^{\text{th}}$  roots of unity is exactly  $s_0 = d_1 \dots d_{\mathcal{N}-1} d_{\mathcal{N}}$ , we have all the solutions.

We thus have  $r_0 + s_0 = 1 + d_1 \dots d_{\mathcal{N}-1} d_{\mathcal{N}}$  sheets. In one of them we have  $z_{\mathcal{N}} \sim O(z_0^{r_{\mathcal{N}}})$ , and in the  $d_1 \dots d_{\mathcal{N}-1} d_{\mathcal{N}}$  others we have  $z_0 \sim O(z_{\mathcal{N}}^{s_0})$ .

## 4.2 Algebraic geometry

We now consider the algebraic curve  $\mathcal{E}$  in a more geometric language. An abstract point  $p \in \mathcal{E}$  can be represented as a couple  $(z_0, z_1)$  such that  $z_1 = Z_1(z_0)$ , or by any other parametrization. For instance it can be described as a point of  $\mathbf{C}^{\mathcal{N}}$ , at the intersection of  $\mathcal{N}$  codimension 1 manifolds.

Algebraic geometry is an active and important part of mathematics, and lots of tools have been invented to describe the geometry of algebraic Riemann surfaces. We refer the reader to [16, 17] for an introduction.

A Riemann surface is locally homomorphic to the complex plane  $\mathbf{C}$ , that means that small domains of  $\mathcal{E}$  can be mapped on small domains of  $\mathbf{C}$ . The map, which is one-to-one and holomorphic in that domain is called a local parameter, let us call it:

$$p \in \mathcal{E} \rightarrow x(p) \in \mathbf{C} \quad (58)$$

Any complex valued analytic function on  $\mathcal{E}$  can be locally represented by an analytic function of  $x(p)$ .

For instance,  $z_0 = z_0(p)$  is often a good local parameter on  $\mathcal{E}$ , except when  $z_0$  approaches a singularity. Therefore all functions on  $\mathcal{E}$  can be locally written as functions of  $z_0$ . They can also be locally written as functions of any  $z_k$ .

Notice that the function  $z_0(p)$  is not injective, it takes the same value  $z_0$  for different points  $p$ , namely the same value of  $z_0$  corresponds to  $1 + s_0$

points  $p$ . Therefore the function  $z_1(p)$  which is a well defined monovalued function on  $\mathcal{E}$  is a multivalued function of  $z_0$ . The  $z_0$ -sheets are domains of  $\mathcal{E}$  where the function  $z_0(p)$  is injective. In particular in each sheet, there is only one point  $p_\infty$  where  $z_0 \rightarrow \infty$ .

Let  $p_{\infty+}$  be the point in the physical sheet, such that  $z_0(p_{\infty+}) = \infty$ . In the vicinity of  $p_{\infty+}$ , we have eq.(54):

$$z_k(p) \underset{p \rightarrow p_{\infty+}}{\sim} z_0^{r_k}(p) \quad (59)$$

i.e.  $Z_k(z_0)$  is analytical, which indicates that  $z_0$  is a good local parameter near  $p_{\infty+}$ .

Now, let  $p_{\infty-}$  be the point at  $\infty$  in the  $z_{\mathcal{N}}$ -physical sheet, i.e. where  $z_{\mathcal{N}}(p_{\infty-}) = \infty$  and with behavior eq.(56):

$$z_k(p) \underset{p \rightarrow p_{\infty-}}{\sim} z_{\mathcal{N}}^{s_k}(p) \quad (60)$$

i.e.  $Z_k(z_{\mathcal{N}})$  is analytical, which indicates that  $z_{\mathcal{N}}$  is a good local parameter near  $p_{\infty-}$ . In the vicinity of  $p_{\infty-}$ , the function  $z_0(p)$  takes the same value  $z_0$ ,  $s_0$  times, therefore there are  $s_0$  sheets which meet at  $p_{\infty-}$ . It is clear that  $z_0$  is not a local parameter near  $p_{\infty-}$ , but  $z_{\mathcal{N}}$  is.

It is also clear that there can be no other point  $p$  such that  $z_0(p) = \infty$ . The algebraic curve  $\mathcal{E}$  has only two points at  $\infty$ .

Notice that the intermediate  $z_k$ 's with  $0 < k < \mathcal{N}$  are not appropriate local coordinates near  $p_{\infty+}$  neither near  $p_{\infty-}$  (unless many of the  $d_k$ 's are equal to 1).

Let us summarize as follows:

- The function  $z_k(p)$  has a pole of degree  $r_k = d_0 d_1 \dots d_{k-1}$  near  $p_{\infty+}$ .  $z_0(p)$  is a local parameter near  $p_{\infty+}$ .
- The function  $z_k(p)$  has a pole of degree  $s_k = d_{k+1} \dots d_{\mathcal{N}-1} d_{\mathcal{N}}$  near  $p_{\infty-}$ .  $z_{\mathcal{N}}(p)$  is a local parameter near  $p_{\infty-}$ .
- The function  $z_k(p)$  has no other pole.

**Remark:** From eq.(53), it is easy to prove by recursion that:

$$\text{Res}_{p_{\infty+}} z_{k-1} dz_k = -T = -\text{Res}_{p_{\infty+}} z_{k+1} dz_k, \quad \text{Res}_{p_{\infty-}} z_{k+1} dz_k = -T = -\text{Res}_{p_{\infty-}} z_{k-1} dz_k \quad (61)$$



### 4.2.1 Genus and cycles

Let  $g$  be the genus of  $\mathcal{E}$ . It can be proved (using the Riemann-Roch theorem, or the method of [19]) that

$$g \leq g_{\max} := \prod_{k=0}^{\mathcal{N}} d_k - 1 \quad (62)$$

Notice that the polynomial  $P$  has  $1 + g_{\max}$  coefficients, and its leading coefficient is fixed, i.e.  $g_{\max}$  is the number of coefficients of  $P$  not fixed by the loop equations.

Let  $\mathcal{A}_i, \mathcal{B}_i (i = 1, \dots, g)$  be a canonical basis of non-trivial cycles on  $\mathcal{E}$ :

$$\mathcal{A}_i \cap \mathcal{B}_j = \delta_{ij} \quad (63)$$

the choice of non-trivial cycles is not unique, and we will see below one possible convenient choice.

### 4.2.2 Endpoints

If  $z_l$  is considered as a function of  $z_k$  ( $k \neq l$ ), it has singularities (branch points) everytime that:

$$dz_k(p) = 0 \quad (64)$$

Indeed at such a point we have:  $z'_l(z_k) = dz_l/dz_k \rightarrow \infty$ .

The zeroes of  $dz_k(p)$  are called the endpoints, and are noted:

$$e_{k,i} , \quad (i = 1, \dots, r_k + s_k + 2g) \quad (65)$$

Generically, the zeroes of  $dz_k$  are simple zeroes and they are all distinct, which means that  $z_k$  behaves as a quadratic function of a local parameter  $x$ , while  $z_l$  behaves linearly in  $x$ . Therefore,  $z_l(z_k)$  has a square root branch point near  $z_k(e_{k,i})$ .

### 4.2.3 Critical points

It may happen that some endpoints coincide. This is called a critical point. It is not a generic situation, it happens only if the potentials  $V_k$ 's are fine tuned to some critical potentials. The critical points are relevant for finding the representations of  $(p, q)$  minimal conformal models. A  $(p, q)$  critical point, is

such that there exist  $k \neq l$ , and a point  $e \in \mathcal{E}$  such that:  $dz_k$  has a zero of degree  $p - 1$  at  $e$  and  $dz_l$  has a zero of degree  $q - 1$  at  $e$ . Near  $e$ ,  $z_k$  as a function of  $z_l$  behaves with a  $p/q$  exponent.

From now on, we assume that the potentials  $V_k$  are not critical, i.e. the endpoints are all simple and distinct.

#### 4.2.4 Cuts

The cuts in  $\mathcal{E}$  are the contours which border the sheets. Viewed in the  $z_k$ -plane they are lines in  $\mathbf{C}$  joining two endpoints.

That choice of lines (i.e. the choice of sheets) is largely arbitrary, provided that they indeed go through the endpoints.

There is a canonical choice, which comes from [13], given by the contours where the asymptotics of the biorthogonal polynomials and their Fourier and Hilbert transforms, have discontinuities, i.e. some Stokes lines. That canonical choice is defined as follows: the set of  $z_k$ -cuts (contours on  $\mathcal{E}$ ) is the set of points  $p \in \mathcal{E}$  such that there exists  $p' \neq p$  with  $z_k(p') = z_k(p)$  and

$$\Re \left( \int_p^{p'} Z_{k+1} dz_k \right) = 0 \quad (66)$$

remark that for all  $p$  and  $p'$  such that  $z_k(p) = z_k(p')$  we have:

$$\int_p^{p'} z_{k+1} dz_k = \int_p^{p'} (V'_k(z_k) - z_{k-1}) dz_k = - \int_p^{p'} z_{k-1} dz_k \quad (67)$$

so that the canonical cuts are left unchanged if we reverse the order of the chain.

### 4.3 Sheet geometry

We are going to describe briefly the sheet geometry, see fig.(4.3) for a better vizualization.

The equation  $z_k(p) = z_k$  has  $r_k + s_k$  solutions for  $p$ , i.e. the curve  $\mathcal{E}$  is divided into  $r_k + s_k$  sheets. When  $z_k \rightarrow \infty$ ,  $r_k$  solutions approach  $p_{\infty+}$ , while  $s_k$  approach  $p_{\infty-}$ . In other words,  $r_k$  sheets contain  $p_{\infty+}$ , while  $s_k$  sheets contain  $p_{\infty-}$ .

Let us denote  $\mathcal{C}_k$  the contour which separates the reunion of sheets containing  $p_{\infty+}$  from the reunion of sheets containing  $p_{\infty-}$ , oriented such that

it turns around  $p_{\infty+}$  in the positive direction. That contour  $\mathcal{C}_k$  will play an important role later. Note that  $\mathcal{C}_k$  is not necessarily connected (as in our example fig. 4.3).

Let us also denote:

- $p_{+j,k}(z_k)$  ( $j = 1, \dots, r_k$ ) all the solutions of  $z_k(p) = z_k$  which are on the same side of  $\mathcal{C}_k$  than  $p_{\infty+}$ ,
- $p_{-j,k}(z_k)$  ( $j = 1, \dots, s_k$ ) all the solutions of  $z_k(p) = z_k$  which are on the same side of  $\mathcal{C}_k$  than  $p_{\infty-}$ .

#### 4.4 Filling fractions

It is well known that the discontinuities of  $W_0(z_0)$ , i.e. the discontinuities of  $Z_{-1}(z_0)$  along the  $Z_0$ -cuts in the physical sheet are related to the large  $N$  average density of eigenvalues of the matrix  $M_0$ :

$$\rho_0(z_0) = -\frac{1}{2i\pi}(W_0(z_0 + i0) - W_0(z_0 - i0)) \quad (68)$$

The support of  $\rho_0$  is the set of cuts of  $Z_{-1}(z_0)$ , with endpoints the  $z_0(e_{0,i})$ 's which belong to the physical sheet.

The ratio of the average number of eigenvalues in a given connected component of the support of  $\rho_0$  to the total number  $N$  of eigenvalues is a contour integral:

$$\frac{1}{T}\epsilon_i = \int_{z_0(e_{0,i})}^{z_0(e_{0,i+1})} \rho(z_0) dz_0 = -\frac{1}{2i\pi T} \oint z_{-1} dz_0 \quad (69)$$

where the contour of integration on  $\mathcal{E}$  is one of the  $z_0$ -cuts defined in section 4.2.4, oriented in the clockwise direction.  $\epsilon_i$  is non-zero only if the contour is a non-trivial contour, or if it encloses a pole  $p_{\infty+}$  or  $p_{\infty-}$ .

Therefore, there are at most  $g + 1$  possible values of  $\epsilon_i$ , the support of the density has at most  $g + 1$  connected components.

It is possible to choose the potentials  $V_1, \dots, V_N$  such that all the non trivial cycles  $\mathcal{A}_i$  are in the physical sheet, and such that the cycles  $\mathcal{A}_i$  are all cuts. The filling fractions in that case are the  $\mathcal{A}$ -cycles integrals:

$$\epsilon_i := -\frac{1}{2i\pi} \oint_{\mathcal{A}_i} z_{-1} dz_0 \quad (70)$$

For more generic potentials, we define the filling fractions by eq.(70), eventhough they don't really correspond to numbers of eigenvalues. The cuts are integer homological linear combination of  $\mathcal{A}$ -cycles.

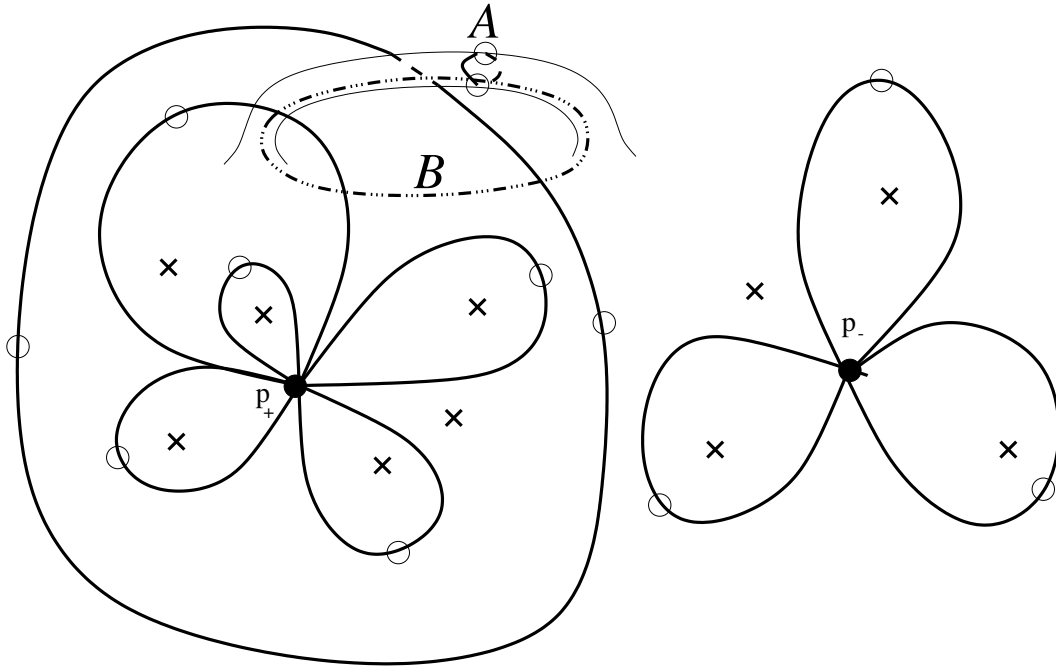


Figure 1: Sheets for  $z_k$ : example with  $r_k = 6$ ,  $s_k = 4$  and  $g = 1$ . The two black circles represent  $p_{\infty+}$  and  $p_{\infty-}$ . The empty circles are the  $r_k + s_k + 2g$  endpoints. The thick lines, going through the endpoints, are the cuts separating the sheets. Each region separated by lines is a sheet, there are  $r_k + s_k$  sheets. The sheets form “flowers” near  $p_{\infty+}$  and  $p_{\infty-}$  (the petals have an angle  $2\pi/r_k$  near  $p_{\infty+}$  and  $2\pi/s_k$  near  $p_{\infty-}$ ). There are  $g$  “handles”, with  $\mathcal{A}$  and  $\mathcal{B}$  cycles: here  $\mathcal{A}$  can be chosen as a cut, and  $\mathcal{B}$  is represented as a thick dash-dot line. Each cross  $\mathbf{x}$  (one in each sheet) represents a point  $p \in \mathcal{E}$ , such that  $z_k(p) = z_k$ .

We also define the chemical potentials as the  $\mathcal{B}$ -cycle integrals:

$$\Gamma_i := - \oint_{\mathcal{B}_i} z_{-1} dz_0 \quad (71)$$

We will see in part 6.1 that this description remains valid for the densities of all matrices  $M_k$  (and not only  $M_0$ ). We will prove (see eq.(160)) that  $\rho_k(z_k)$  the density of eigenvalues of  $M_k$  is the discontinuity of  $z_{k+1}$  (or equivalently  $z_{k-1}$ ) along the cuts.

It is easy to prove, by integration by parts, that:

$$\oint_{\mathcal{A}_i} z_{-1} dz_0 = \oint_{\mathcal{A}_i} z_{k-1} dz_k = - \oint_{\mathcal{A}_i} z_{k+1} dz_k \quad (72)$$

$$\oint_{\mathcal{B}_i} z_{-1} dz_0 = \oint_{\mathcal{B}_i} z_{k-1} dz_k = - \oint_{\mathcal{B}_i} z_{k+1} dz_k \quad (73)$$

If we anticipate on section 6.1, this proves that the filling fractions are the same for all densities, up to integer linear combinations. If the support of the density of eigenvalues of  $M_k$  has  $m_k \leq g+1$  connected components  $[a_{k,i}, b_{k,i}]$  ( $i = 1, \dots, m_k$ ), the filling fractions in each connected component are:

$$\int_{a_{k,i}}^{b_{k,i}} \rho_k(z_k) dz_k = \frac{1}{T} \sum_{j=1}^{g+1} A_{k,i,j} \epsilon_j \quad (74)$$

where the coefficients  $A_{k,i,j}$  are integers (possibly nul or negative).

## 4.5 $\mathcal{A}$ or $\mathcal{B}$ cycles fixed ?

As we said before, the loop equation is not sufficient to determine all the unknown coefficients of  $E$ . Some additional hypothesis are needed, two of them are often considered in the litterature:

- **condition B:** if all the potentials are bounded from below, and the partition function is well defined, one is interested in finding large  $N$  limits for the free energy and various expectation values of traces of powers of the matrices. In that approach, one has to find a solution of the loop equation which gives an absolute minimum of the free energy. The genus  $g$  and the filling fractions  $\epsilon_i$  are not known, they are

determined by the minimization condition  $\partial F/\partial \epsilon_i = 0$ , which reduces to:

$$\forall i = 1, \dots, g \quad \oint_{\mathcal{B}_i} z_1 dz_0 = 0 \quad (75)$$

All the  $\mathcal{B}$  cycle integrals must vanish. That condition is sufficient to determine  $g$  and all the  $\epsilon_i$ , and it is sufficient to determine the unknown coefficients of  $E$ .

This allows to find the large  $N$  limit of the free energy. Subleading large  $N$  corrections, i.e. the so-called topological large  $N$  expansion of the free energy, exist only if  $g = 0$ , as was shown in [6, 10]. If the genus  $g$  is  $\geq 1$ , there exist asymptotics for the subleading corrections to the free energy, which have no  $1/N$  expansion, they are oscillating functions.

- **condition A:** if one is only interested in the formal large  $N$  expansion of the free energy, one has to find a solution of the loop equation which corresponds to the perturbation of a given local extremum of the free energy. Therefore,  $g$  and the  $\epsilon_i$ 's are fixed parameters which characterize the minimum around which we perform the perturbative expansion (moduli). The following equations

$$\forall i = 1, \dots, g \quad \oint_{\mathcal{A}_i} z_1 dz_0 = 2i\pi\epsilon_i \quad (76)$$

are sufficient to determine all the unknown coefficients of  $E$ .

In the following, we will always assume that we have condition A, unless specified. Most of the results obtained for condition A, immediately translate to condition B, by exchanging the roles of  $\mathcal{A}$  and  $\mathcal{B}$ -cycles, like the large  $N$  free energy. But some quantities, like the large  $N$  2-loop functions, get contributions from the oscillating asymptotics, and will not be computed in this article. So, from now on, we assume condition A.

## 4.6 Parametrization

It is possible to parametrize the algebraic curve  $\mathcal{E}$  in terms of  $\theta$  functions.

Given  $g$  and a canonical basis of non-trivial cycles  $\mathcal{A}_i, \mathcal{B}_i (i = 1, \dots, g)$ , it is known [16, 17] that there exists a unique basis of holomorphic 1-forms  $du_i$

on  $\mathcal{E}$  such that:

$$\oint_{\mathcal{A}_i} du_j = \delta_{ij} \quad (77)$$

The matrix of  $\mathcal{B}$  periods is defined as:

$$\tau_{ij} := \oint_{\mathcal{B}_i} du_j = \tau_{ji} \quad (78)$$

$\tau_{ij}$  is a symmetric  $g \times g$  matrix, with positive imaginary part.

If we choose an arbitrary basepoint  $p_0 \in \mathcal{E}$ , we define the Abel map:

$$\vec{u}(p) := \int_{p_0}^p d\vec{u} \quad (79)$$

which defines an embedding of  $\mathcal{E}$  into  $\mathbf{C}^g$ .

For any  $0 \leq k \leq \mathcal{N}$ ,  $z_k(p)$  is a function on  $\mathcal{E}$ , with a pole of degree  $r_k$  at  $p_{\infty+}$ , and a pole of degree  $s_k$  at  $p_{\infty-}$ .  $z_k$  must have  $r_k + s_k$  zeroes, which we denote  $p_{k,i}(0)$ , ( $i = 1, \dots, r_k + s_k$ ). The zeroes must satisfy:

$$\sum_{i=1}^{r_k+s_k} \vec{u}(p_{k,i}(0)) = r_k \vec{u}(p_{\infty+}) + s_k \vec{u}(p_{\infty-}) \quad (80)$$

then:

$$z_k(p) = A_k \frac{\prod_{i=1}^{r_k+s_k} \theta(\vec{u}(p) - \vec{u}(p_{k,i}(0)) - \vec{z}; \tau)}{\theta(\vec{u}(p) - \vec{u}(p_{\infty+}) - \vec{z}; \tau)^{r_k} \theta(\vec{u}(p) - \vec{u}(p_{\infty-}) - \vec{z}; \tau)^{s_k}} \quad (81)$$

where  $\theta$  is the Riemann theta function [16,17],  $\vec{z}$  is an arbitrary non-singular zero of  $\theta$ . The ratio of the RHS and the LHS is a function on  $\mathcal{E}$  with no pole, therefore it must be a constant, which we call  $A_k$ .

We have a parametrization of  $\mathcal{E}$  in terms of  $\theta$  functions.

## 4.7 Two-loop functions and Bergmann kernel

### 4.7.1 The Bergmann kernel

The Bergmann kernel  $B(p; p')$  is the unique bilinear differential form on  $\mathcal{E} \times \mathcal{E}$  with the following properties:

- $B(p; p')$ , as a function of  $p$ , is a meromorphic form, with only one double pole at  $p = p'$  with no residue, and such that in any local parameter  $z(p)$  we have:

$$B(p; p') \underset{p \rightarrow p'}{\sim} \frac{dz(p) dz(p')}{(z(p) - z(p'))^2} \quad (82)$$

- $\forall i = 1, \dots, g \quad \oint_{p \in \mathcal{A}_i} B(p, p') = 0.$

#### 4.7.2 The 2-point function

For  $0 \leq k \leq \mathcal{N}$  and  $0 \leq l \leq \mathcal{N}$ , define the following meromorphic differential forms on  $\mathcal{E} \times \mathbf{C}$ :

$$B_{k,l}(p; z) := \left. \frac{\partial z_{k-1}(p)}{\partial V_l(z)} \right|_{z_k(p)} dz_k(p) dz \quad (83)$$

$$\tilde{B}_{k,l}(p; z) := \left. \frac{\partial z_{k+1}(p)}{\partial V_l(z)} \right|_{z_k(p)} dz_k(p) dz \quad (84)$$

We have the “thermodynamic identity”:

$$\tilde{B}_{k,l}(p; z) = -B_{k+1,l}(p; z) \quad (85)$$

and  $z_{k-1} + z_{k+1} = V'_k(z_k)$  implies:

$$B_{k,l}(p; z) + \tilde{B}_{k,l}(p; z) = -\delta_{k,l} \frac{dz_l(p) dz}{(z - z_l(p))^2} \quad (86)$$

We are going to prove below that:

$$\begin{aligned} B_{k,l}(p; z) &= -\sum_{j=1}^{s_l} B(p; p_{-j,l}(z)) & 0 \leq k \leq l \\ B_{k,l}(p; z) &= \sum_{j=1}^{r_l} B(p; p_{+j,l}(z)) & l < k \leq \mathcal{N} \\ \tilde{B}_{k,l}(p; z) &= \sum_{j=1}^{s_l} B(p; p_{-j,l}(z)) & 0 \leq k < l \\ \tilde{B}_{k,l}(p; z) &= -\sum_{j=1}^{r_l} B(p; p_{+j,l}(z)) & l \leq k \leq \mathcal{N} \end{aligned} \quad (87)$$

where  $B$  is the Bergmann kernel, and the  $p_{j,l}$ ’s have been defined in 4.3.

#### Proof of eq.(87)

First notice that eq.(85) and eq.(86) imply:

$$\begin{aligned} B_{k,l}(p; z) &= B_{0,l}(p; z) & 0 \leq k \leq l \\ B_{k,l}(p; z) &= -\tilde{B}_{\mathcal{N},l}(p; z) & l < k \leq \mathcal{N} \end{aligned} \quad (88)$$



and:

$$B_{0,l}(p; z) + \tilde{B}_{\mathcal{N},l}(p; z) = -\frac{dz_l(p) dz}{(z - z_l(p))^2} \quad (89)$$

Consider the two-loop functions introduced in eq.(27):

$$W_{0,l}(z; z') = \left\langle \text{tr} \frac{1}{z - M_0} \text{tr} \frac{1}{z' - M_l} \right\rangle_{\text{conn}} = \frac{B_{0,l}(p_{+1,0}(z); z')}{dz_0(p_{+1,0}(z))dz'} \quad (90)$$

$$W_{\mathcal{N},l}(z; z') = \left\langle \text{tr} \frac{1}{z - M_{\mathcal{N}}} \text{tr} \frac{1}{z' - M_l} \right\rangle_{\text{conn}} = \frac{\tilde{B}_{\mathcal{N},l}(p_{-1,\mathcal{N}}(z); z')}{dz_{\mathcal{N}}(p_{-1,\mathcal{N}}(z))dz'} \quad (91)$$

The expectation values  $\left\langle \text{tr} \frac{1}{z - M_0} \text{tr} \frac{1}{z' - M_l} \right\rangle$  and  $\left\langle \text{tr} \frac{1}{z - M_{\mathcal{N}}} \text{tr} \frac{1}{z' - M_l} \right\rangle$  are well defined integrals for  $z$  in the physical sheet, therefore  $B_{0,l}(p; z')$  can have no pole when  $z' = z_l(p)$  if  $p$  is in the  $z_0$ -physical sheet, and  $\tilde{B}_{\mathcal{N},l}(p; z')$  can have no pole when  $z' = z_l(p)$  if  $p$  is in the  $z_{\mathcal{N}}$ -physical sheet. Note that the derivatives  $\partial/\partial V_l(z')$  are formally defined only for large  $z'$ , which implies that:

- $B_{0,l}(p; z')$  can have no pole when  $p = p_{+j,l}(z')$ ,  $j = 1, \dots, r_l$ ,
- $\tilde{B}_{\mathcal{N},l}(p; z')$  can have no pole when  $p = p_{-j,l}(z')$ ,  $j = 1, \dots, s_l$ .

Since the RHS of eq.(89) has poles at all  $p = p_{\pm j,l}(z')$ , we must have:

- $B_{0,l}(p; z')$  can have no pole when  $p = p_{+j,l}(z')$ ,  $j = 1, \dots, r_l$ ,
- $B_{0,l}(p; z')$  has double poles when  $p = p_{-j,l}(z')$ ,  $j = 1, \dots, s_l$ .

That implies that:

$$\overline{B}_l(p; z') := B_{0,l}(p; z') + \sum_{j=1}^{s_l} B(p; p_{-j,l}(z')) \quad (92)$$

has no pole when  $z_l(p) = z'$ .

It obeys the following properties:

- Since  $W_0(z_0)$  behaves as  $1/z_0$  near  $\infty$  in the physical sheet,  $W_{0,l}(z_0; z')$  must behave as  $O(1/z_0^2)$  near  $p_{\infty+}$ , while  $dz_0$  has a double pole at  $p_{\infty+}$ . Therefore  $B_{0,l}(p; z')$  has no pole when  $p \rightarrow p_{\infty+}$ , and thus  $\overline{B}_l(p; z')$  has no pole when  $p \rightarrow p_{\infty+}$ .
- Similarly,  $\tilde{B}_{\mathcal{N},l}(p; z')$  has no pole when  $p \rightarrow p_{\infty-}$ , and with eq.(89), that implies that  $B_{0,l}(p; z')$  has no pole when  $p \rightarrow p_{\infty-}$ , and thus  $\overline{B}_l(p; z')$  has no pole when  $p \rightarrow p_{\infty-}$ .

- Near an endpoint  $e_{k,i}$ ,  $Z_{k-1}(z_k)$  has a square root singularity, i.e. the derivative  $\partial z_{k-1}(p)/\partial V_l(z')|_{z_k(p)}$  has an inverse square root singularity, i.e. a simple pole which is exactly compensated by the zero of  $dz_k(p)$ . Therefore  $B_{k,l}(p; z')$  has no pole when  $p \rightarrow e_{k,i}$ , and so for  $\overline{B}_l(p; z')$ .
- Near any other point,  $z_{k-1}$  is an analytical function of  $z_k$ , and thus  $B_{k,l}$  is analytical.
- $B_{k,l}$  must satisfy:

$$\oint_{p \in \mathcal{A}_i} B_{k,l}(p; z') = 0 \quad (93)$$

and so must  $\overline{B}_l$ .

Finally we find that  $\overline{B}_l$  is a meromorphic form with no pole, with all its  $\mathcal{A}$  cycle integrals vanishing, therefore:  $\overline{B}_l = 0$ , QED.

**Remark:**

We have:

$$B(p, p') = d_p d_{p'} \ln \theta(\vec{u}(p) - \vec{u}(p') - \vec{z}) \quad (94)$$

Notice that

$$du_i(p') = \frac{1}{2i\pi} \oint_{p \in \mathcal{B}_i} B(p, p') = dz_0(p') \frac{1}{2i\pi} \frac{\partial \Gamma_i}{\partial V_0(z_0(p'))} \quad (95)$$

Once  $B$  is known, eq.(95) give an explicit way of computing the Abel map  $\vec{u}(p)$ .

## 4.8 Abelian differential of the 3rd kind and temperature

### 4.8.1 Abelian differential of the 3rd kind

There is a unique abelian differential  $dS$  on  $\mathcal{E}$  with only two simple poles at  $p_{\infty\pm}$ , such that:

$$\text{Res}_{p_{\infty+}} dS = -\text{Res}_{p_{\infty-}} dS = 1, \quad \forall i = 1, \dots, g \quad \oint_{\mathcal{A}_i} dS = 0 \quad (96)$$

It has the property that:

$$dS(p) = \int_{p'=p_{\infty-}}^{p_{\infty+}} B(p, p') \quad (97)$$

We define:

$$\eta_i := \oint_{\mathcal{B}_i} dS \quad (98)$$

And, given a basepoint  $p_0$ , we define the following multivalued function on  $\mathcal{E}$ :

$$\Lambda(p) := e^{-\int_{p_0}^p dS} \quad (99)$$

$\Lambda$  has a simple pole at  $p_{\infty+}$  and a single zero at  $p_{\infty-}$ , therefore the following quantities are well defined:

$$\gamma := \lim_{p \rightarrow p_{\infty+}} \frac{z_0(p)}{\Lambda(p)}, \quad \tilde{\gamma} := \lim_{p \rightarrow p_{\infty-}} z_{\mathcal{N}}(p) \Lambda(p) \quad (100)$$

Notice that the product  $\gamma\tilde{\gamma}$  is independent of the choice of the basepoint  $p_0$ .

#### 4.8.2 Derivatives with respect to $T$

Consider the abelian differentials:

$$dS_k(p) := \left. \frac{\partial z_{k-1}(p)}{\partial T} \right|_{z_k(p)} dz_k(p), \quad d\tilde{S}_k(p) := \left. \frac{\partial z_{k+1}(p)}{\partial T} \right|_{z_k(p)} dz_k(p) \quad (101)$$

We clearly have (from  $z_{k+1} + z_{k-1} = V'_k(z_k)$  and thermodynamic identity):

$$dS_k(p) = -d\tilde{S}_k(p) = dS_{k+1}(p) \quad (102)$$

Therefore  $dS_k$  and  $d\tilde{S}_k$  are independent of  $k$ .

They have the following properties:

- near  $p_{\infty+}$ ,  $z_{-1} \sim T/z_0 + O(z_0^{-2})$ , thus  $dS_0 \sim dz_0/z_0$ , i.e.  $dS_0$  has a single pole at  $p_{\infty+}$ , with residue  $-1$ .
- Similarly  $d\tilde{S}_{\mathcal{N}}$  has a single pole at  $p_{\infty-}$ , with residue  $-1$ , i.e.  $dS_0$  has a single pole at  $p_{\infty-}$ , with residue  $+1$ .
- near an endpoint  $e_{k,i}$ ,  $Z_{k-1}(z_k)$  has a square root singularity, thus  $\left. \frac{\partial z_{k-1}}{\partial T} \right|_{z_k}$  has an inverse square root singularity, i.e. a single pole at  $e_{k,i}$ , which is compensated by the single zero of  $dz_K$ , and therefore  $dS_0$  has no pole at  $e_{k,i}$ .
- near any other point,  $z_k$  is a local parameter and  $z_{k+1}$  is analytical in  $z_k$ , i.e.  $dS_0$  has no pole.

- The  $\mathcal{A}_i$  cycle integrals of  $dS_0$  vanish.

There is a unique abelian differential with such properties, it is the abelian differential of the third kind defined in eq.(96):

$$d\tilde{S}_k(p) = -dS_k(p) = dS(p) \quad (103)$$

In particular, we have that:

$$\frac{\partial}{\partial T}\Gamma_i = \eta_i, \quad \frac{\partial}{\partial T}\epsilon_i = 0 \quad (104)$$

#### 4.9 Derivatives with respect to $\epsilon_i$ (condition A)

Consider:

$$du_{k,i}(p) := \frac{1}{2i\pi} \left. \frac{\partial z_{k-1}(p)}{\partial \epsilon_i} \right|_{z_k(p)} dz_k(p), \quad d\tilde{u}_{k,i}(p) := \frac{1}{2i\pi} \left. \frac{\partial z_{k+1}(p)}{\partial \epsilon_i} \right|_{z_k(p)} dz_k(p) \quad (105)$$

We clearly have (from  $z_{k+1} + z_{k-1} = V'_k(z_k)$  and thermodynamic identity):

$$du_{k,i}(p) = -d\tilde{u}_{k,i}(p) = du_{k+1,i}(p) \quad (106)$$

Therefore  $du_{k,i}$  and  $d\tilde{u}_{k,i}$  are independent of  $k$ .

Following the same lines as in the previous section, we show that  $du_{k,i}$  and  $d\tilde{u}_{k,i}$  have no poles, i.e. they are holomorphic one-forms. Moreover we have:

$$\oint_{\mathcal{A}_j} d\tilde{u}_{k,i} = \frac{\partial \epsilon_j}{\partial \epsilon_i} = \delta_{i,j} \quad (107)$$

There exists a unique set of holomorphic one-forms with those properties, it is the holomorphic forms  $du_i$  introduced in eq.(77), therefore:

$$d\tilde{u}_{k,i} = -du_{k,i} = du_i \quad (108)$$

This shows that:

$$\frac{\partial \Gamma_j}{\partial \epsilon_i} = 2i\pi \tau_{i,j} \quad (109)$$

where  $\tau_{i,j}$  is the matrix of periods introduced in eq.(78).

## 5 Large $N$ free energy

### 5.1 The large $N$ free energy

The free energy  $F$  defined in eq.(2) has a large  $N$  limit:

$$F = F^{(0)} + O(N^{-2}) \quad (110)$$

We prove below, a generalization of the formula of [3]:

$$\begin{aligned} 2F^{(0)} &= \sum_{k=0}^{\mathcal{N}} \operatorname{Res}_{p_{\infty+}} (V_k(z_k) - \frac{1}{2} z_k V'_k(z_k)) z_{k+1} dz_k \\ &\quad + T\mu + \sum_i \epsilon_i \Gamma_i - T^2(1 + \mathcal{N} \ln T) \\ &= \sum_{k=0}^{\mathcal{N}} \operatorname{Res}_{p_{\infty-}} (V_k(z_k) - \frac{1}{2} z_k V'_k(z_k)) z_{k-1} dz_k \\ &\quad + T\mu + \sum_i \epsilon_i \Gamma_i - T^2(1 + \mathcal{N} \ln T) \end{aligned} \quad (111)$$

where  $\mu$  is the generalized version of [3] and is defined as follows for any  $p \in \mathcal{E}$ :

$$\begin{aligned} \mu &:= \int_{p_{\infty+}}^p (\frac{T}{z_0} - z_{-1}) dz_0 + \int_{p_{\infty-}}^p (\frac{T}{z_{\mathcal{N}}} - z_{\mathcal{N}+1}) dz_{\mathcal{N}} - T \ln z_0(p) - T \ln z_{\mathcal{N}}(p) \\ &\quad + \sum_{k=0}^{\mathcal{N}} V_k(z_k(p)) - \sum_{k=1}^{\mathcal{N}} z_{k-1}(p) z_k(p) \end{aligned} \quad (112)$$

$\mu$  is independent of  $p$  (indeed  $d\mu$  is a telescopic sum which cancels completely).

**Proof of eq.(111):** We remind that we assume condition A. We define:

$$\begin{aligned} 4K &:= 2T\mu + 2 \sum_i \epsilon_i \Gamma_i - 4F^{(0)} + \sum_{k=0}^{\mathcal{N}} \operatorname{Res}_{p_{\infty+}} (V_k(z_k) - \frac{1}{2} z_k z_{k+1}) z_{k+1} dz_k \\ &\quad + \sum_{k=0}^{\mathcal{N}} \operatorname{Res}_{p_{\infty-}} (V_k(z_k) - \frac{1}{2} z_k z_{k-1}) z_{k-1} dz_k \end{aligned} \quad (113)$$

Notice that  $\operatorname{Res}_{p_{\infty+}} z_k z_{k-1} z_{k+1} dz_k + \operatorname{Res}_{p_{\infty-}} z_k z_{k-1} z_{k+1} dz_k = 0$ , so that the expression in eq.(113) is the same as in eq.(111).

Let us compute:

$$\begin{aligned}
4 \frac{\partial K}{\partial V_l(z)} &= \sum_k \operatorname{Res}_{p_{\infty+}} \frac{\partial V_k(z_k)}{\partial V_l(z)} z_{k+1} dz_k + \sum_k \operatorname{Res}_{p_{\infty-}} \frac{\partial V_k(z_k)}{\partial V_l(z)} z_{k-1} dz_k + 4TW_l(z) \\
&+ \sum_k \operatorname{Res}_{p_{\infty+}} \left. \frac{\partial z_{k+1}}{\partial V_l(z)} \right|_{z_k} (V_k(z_k) - z_k z_{k+1}) dz_k \\
&+ \sum_k \operatorname{Res}_{p_{\infty-}} \left. \frac{\partial z_{k-1}}{\partial V_l(z)} \right|_{z_k} (V_k(z_k) - z_k z_{k-1}) dz_k \\
&+ 2T \frac{\partial \mu}{\partial V_l(z)} + 2 \sum_i \epsilon_i \frac{\partial \Gamma_i}{\partial V_l(z)}
\end{aligned} \tag{114}$$

we introduce the following multivalued functions:

$$\zeta_k(p) := \int_{p_{\infty+}}^p \left. \frac{\partial z_{k+1}}{\partial V_l(z)} \right|_{z_k} dz_k, \quad \tilde{\zeta}_k(p) := \int_{p_{\infty-}}^p \left. \frac{\partial z_{k-1}}{\partial V_l(z)} \right|_{z_k} dz_k \tag{115}$$

By integration by parts, and using eq.(20), we have:

$$\begin{aligned}
4 \frac{\partial K}{\partial V_l(z)} &= - \operatorname{Res}_{p_{\infty+}} \frac{z_{l+1}}{z - z_l} dz_l - \operatorname{Res}_{p_{\infty-}} \frac{z_{l-1}}{z - z_l} dz_l + 4TW_l(z) \\
&- \sum_k \operatorname{Res}_{p_{\infty+}} \zeta_k ((V'_k(z_k) - z_{k+1}) dz_k - z_k dz_{k+1}) \\
&- \sum_k \operatorname{Res}_{p_{\infty-}} \tilde{\zeta}_k ((V'_k(z_k) - z_{k-1}) dz_k - z_k dz_{k-1}) \\
&+ 2T \frac{\partial \mu}{\partial V_l(z)} + 2 \sum_i \epsilon_i \frac{\partial \Gamma_i}{\partial V_l(z)} \\
&= - \operatorname{Res}_{p_{\infty+}} \frac{z_{l+1}}{z - z_l} dz_l - \operatorname{Res}_{p_{\infty-}} \frac{z_{l-1}}{z - z_l} dz_l + 4TW_l(z) \\
&- \sum_k \operatorname{Res}_{p_{\infty+}} \zeta_k (z_{k-1} dz_k - z_k dz_{k+1}) \\
&- \sum_k \operatorname{Res}_{p_{\infty-}} \tilde{\zeta}_k (z_{k+1} dz_k - z_k dz_{k-1}) \\
&+ 2T \frac{\partial \mu}{\partial V_l(z)} + 2 \sum_i \epsilon_i \frac{\partial \Gamma_i}{\partial V_l(z)} \\
&= - \operatorname{Res}_{p_{\infty+}} \frac{z_{l+1}}{z - z_l} dz_l - \operatorname{Res}_{p_{\infty-}} \frac{z_{l-1}}{z - z_l} dz_l + 4TW_l(z) \\
&+ \operatorname{Res}_{p_{\infty+}} \zeta_{\mathcal{N}} z_{\mathcal{N}} dz_{\mathcal{N}+1} - \zeta_{-1} z_{-1} dz_0 - \sum_{k=0}^{\mathcal{N}} (\zeta_k - \zeta_{k-1}) z_{k-1} dz_k
\end{aligned}$$

$$\begin{aligned}
& + \operatorname{Res}_{p_{\infty-}} \tilde{\zeta}_0 z_0 dz_{-1} - \tilde{\zeta}_{\mathcal{N}+1} z_{\mathcal{N}+1} dz_{\mathcal{N}} - \sum_{k=0}^{\mathcal{N}} (\tilde{\zeta}_k - \tilde{\zeta}_{k+1}) z_{k+1} dz_k \\
& + 2T \frac{\partial \mu}{\partial V_l(z)} + 2 \sum_i \epsilon_i \frac{\partial \Gamma_i}{\partial V_l(z)}
\end{aligned} \tag{116}$$

Near  $p_{\infty+}$  we have  $z_{-1} \sim T/z_0 + O(z_0^{-2})$ , therefore  $z_0 \sim T/z_{-1} + O(1)$ , and thus  $\zeta_{-1}$  has a zero at  $p_{\infty+}$ . That implies that:

$$\operatorname{Res}_{p_{\infty+}} \zeta_{-1} z_{-1} dz_0 = 0, \quad \operatorname{Res}_{p_{\infty-}} \tilde{\zeta}_{\mathcal{N}+1} z_{\mathcal{N}+1} dz_{\mathcal{N}} = 0 \tag{117}$$

Moreover, notice that  $\left. \frac{\partial z_k}{\partial V_l(z)} \right|_{z_{k-1}} dz_{k-1} = - \left. \frac{\partial z_{k-1}}{\partial V_l(z)} \right|_{z_k} dz_k$ , therefore:

$$\zeta_k(p) - \zeta_{k-1}(p) = \int_{p_{\infty+}}^p \left( \left. \frac{\partial z_{k+1}}{\partial V_l(z)} \right|_{z_k} + \left. \frac{\partial z_{k-1}}{\partial V_l(z)} \right|_{z_k} \right) dz_k = \int_{p_{\infty+}}^p \frac{\partial V'_k(z_k)}{\partial V_l(z)} dz_k = \frac{\delta_{k,l}}{z_l(p) - z} \tag{118}$$

Thus:

$$\begin{aligned}
4 \frac{\partial K}{\partial V_l(z)} &= \operatorname{Res}_{p_{\infty+}} \frac{z_{l-1} - z_{l+1}}{z - z_l} dz_l + \operatorname{Res}_{p_{\infty-}} \frac{z_{l+1} - z_{l-1}}{z - z_l} dz_l + 4TW_l(z) \\
&+ \operatorname{Res}_{p_{\infty+}} \zeta_{\mathcal{N}} z_{\mathcal{N}} dz_{\mathcal{N}+1} + \operatorname{Res}_{p_{\infty-}} \tilde{\zeta}_0 z_0 dz_{-1} + 2T \frac{\partial \mu}{\partial V_l(z)} + 2 \sum_i \epsilon_i \frac{\partial \Gamma_i}{\partial V_l(z)} \\
&= \operatorname{Res}_{p_{\infty+}} \frac{z_{l-1} - z_{l+1}}{z - z_l} dz_l + \operatorname{Res}_{p_{\infty-}} \frac{z_{l+1} - z_{l-1}}{z - z_l} dz_l + 4TW_l(z) \\
&- \operatorname{Res}_{p_{\infty-}} \zeta_{\mathcal{N}} z_{\mathcal{N}} dz_{\mathcal{N}+1} - \operatorname{Res}_{p_{\infty+}} \tilde{\zeta}_0 z_0 dz_{-1} + 2T \frac{\partial \mu}{\partial V_l(z)} + 2 \sum_i \epsilon_i \frac{\partial \Gamma_i}{\partial V_l(z)} \\
&+ \sum_i \frac{1}{2i\pi} \oint_{\mathcal{A}_i} \operatorname{disc}_{\mathcal{A}_i} (\tilde{\zeta}_0 z_0 dz_{-1}) + \sum_i \frac{1}{2i\pi} \oint_{\mathcal{A}_i} \operatorname{disc}_{\mathcal{A}_i} (\zeta_{\mathcal{N}} z_{\mathcal{N}} dz_{\mathcal{N}+1}) \\
&- \sum_i \frac{1}{2i\pi} \oint_{\mathcal{B}_i} \operatorname{disc}_{\mathcal{B}_i} (\tilde{\zeta}_0 z_0 dz_{-1}) - \sum_i \frac{1}{2i\pi} \oint_{\mathcal{B}_i} \operatorname{disc}_{\mathcal{B}_i} (\zeta_{\mathcal{N}} z_{\mathcal{N}} dz_{\mathcal{N}+1})
\end{aligned} \tag{119}$$

where we have used Riemann bilinear identity, and disc means taking discontinuity of the considered multivalued function when crossing cycles.

Now compute  $\partial \mu / \partial V_l(z)$ , all terms cancel but:

$$\frac{\partial \mu}{\partial V_l(z)} = -\frac{1}{z - z_l(p)} - \int_{p_{\infty+}}^p \left. \frac{\partial z_{-1}}{\partial V_l(z)} \right|_{z_0} dz_0 - \int_{p_{\infty-}}^p \left. \frac{\partial z_{\mathcal{N}+1}}{\partial V_l(z)} \right|_{z_{\mathcal{N}}} dz_{\mathcal{N}} \tag{120}$$

which is independent of  $p$ . In particular for  $p = p_{\infty\pm}$ , this proves that  $\tilde{\zeta}_0(p_{\infty+})$  and  $\zeta_{\mathcal{N}}(p_{\infty-})$  are finite, and:

$$\frac{\partial\mu}{\partial V_l(z)} = \tilde{\zeta}_0(p_{\infty+}) = \zeta_{\mathcal{N}}(p_{\infty-}) \quad (121)$$

Thus:

$$\operatorname{Res}_{p_{\infty+}} \tilde{\zeta}_0 z_0 dz_{-1} = T \tilde{\zeta}_0(p_{\infty+}) = T \frac{\partial\mu}{\partial V_l(z)} = T \zeta_{\mathcal{N}}(p_{\infty-}) = \operatorname{Res}_{p_{\infty-}} \zeta_{\mathcal{N}} z_{\mathcal{N}} dz_{\mathcal{N}+1} \quad (122)$$

Now compute the discontinuities (crossing  $\mathcal{B}_i$  is equivalent to going around  $\mathcal{A}_i$ ):

$$\operatorname{disc}_{\mathcal{B}_i} \tilde{\zeta}_0 = \oint_{\mathcal{A}_i} \frac{\partial z_{-1}}{\partial V_l(z)} \Big|_{z_0} dz_0 = -2i\pi \frac{\partial \epsilon_i}{\partial V_l(z)} = 0 \quad (123)$$

this implies that  $\tilde{\zeta}_0$  has no discontinuity along  $\mathcal{B}_i$ , and

$$\operatorname{disc}_{\mathcal{A}_i} \tilde{\zeta}_0 = \oint_{\mathcal{B}_i} \frac{\partial z_{-1}}{\partial V_l(z)} \Big|_{z_0} dz_0 = -\frac{\partial \Gamma_i}{\partial V_l(z)} \quad (124)$$

this implies that  $\tilde{\zeta}_0$  has a constant discontinuity along  $\mathcal{A}_i$ . Moreover  $z_0$  and  $dz_{-1}$  are monovalued, i.e. they have no discontinuity along  $\mathcal{A}_i$  or  $\mathcal{B}_i$ . Thus:

$$\oint_{\mathcal{A}_i} \operatorname{disc}_{\mathcal{A}_i} (\tilde{\zeta}_0 z_0 dz_{-1}) = -\frac{\partial \Gamma_i}{\partial V_l(z)} \oint_{\mathcal{A}_i} z_0 dz_{-1} = -2i\pi \epsilon_i \frac{\partial \Gamma_i}{\partial V_l(z)} \quad (125)$$

Using similar arguments for  $z_{\mathcal{N}}$ , we arrive at:

$$4 \frac{\partial K}{\partial V_l(z)} = \operatorname{Res}_{p_{\infty+}} \frac{z_{l-1} - z_{l+1}}{z - z_l} dz_l + \operatorname{Res}_{p_{\infty-}} \frac{z_{l+1} - z_{l-1}}{z - z_l} dz_l + 4TW_l(z) \quad (126)$$

When  $l = 0$ , by definition eq.(11), we have

$$\begin{aligned} TW_0(z) &= \operatorname{Res}_{p_{\infty+}} \frac{1}{z - z_0} z_1 dz_0 = -\operatorname{Res}_{p_{\infty+}} \frac{1}{z - z_0} z_{-1} dz_0 \\ &= \operatorname{Res}_{p_{\infty-}} \frac{1}{z - z_0} z_{-1} dz_0 = -\operatorname{Res}_{p_{\infty-}} \frac{1}{z - z_0} z_1 dz_0 \end{aligned} \quad (127)$$

therefore:

$$\frac{\partial K}{\partial V_0(z)} = 0 \quad (128)$$



which proves that  $K$  is independent of  $V_0$ . In particular we can choose  $V_0$  quadratic, and then we integrate  $M_0$  out, i.e. we reduce the problem of a chain of lenght  $\mathcal{N}$  with potentials  $V_0, \dots, V_{\mathcal{N}}$  to a chain of lenght  $\mathcal{N} - 1$  with potentials  $V_1 - \frac{z_1^2}{2}, \dots, V_{\mathcal{N}}$ . It is easy to check directly that:

$$F_{\mathcal{N}}(\frac{z_0^2}{2}, V_1, \dots, V_{\mathcal{N}}) = F_{\mathcal{N}-1}(V_1 - \frac{z_1^2}{2}, \dots, V_{\mathcal{N}}) - \frac{T^2}{2} \ln T \quad (129)$$

and thus:

$$K_{\mathcal{N}}(V_0, \dots, V_{\mathcal{N}}) = K_{\mathcal{N}}(\frac{z_0^2}{2}, V_1, \dots, V_{\mathcal{N}}) = K_{\mathcal{N}-1}(V_1 - \frac{z_1^2}{2}, \dots, V_{\mathcal{N}}) + \frac{T^2}{2} \ln T \quad (130)$$

Therefore, by recursion on  $\mathcal{N}$ , we find that  $K$  is independent of all  $V$ 's.  $K$  could still depend on  $T$  and the  $\epsilon_i$ 's.  $K$  has been computed for  $\mathcal{N} = 1$  and  $\mathcal{N} = 0$  [3], and we find:

$$K = \frac{T^2}{2}(1 + \mathcal{N} \ln T) \quad (131)$$

which is the same result as if all potentials are chosen gaussian (see section 8.2).

QED

**Remark:** En route, we have proved that for all  $l$ , the resolvent of the  $l^{\text{th}}$  matrix is:

$$\begin{aligned} W_l(z) &= \frac{1}{T} \text{Res}_{p_{\infty+}} \frac{1}{z - z_l} z_{l+1} dz_l = -\frac{1}{T} \text{Res}_{p_{\infty+}} \frac{1}{z - z_l} z_{l-1} dz_l \\ &= \frac{1}{T} \text{Res}_{p_{\infty-}} \frac{1}{z - z_l} z_{l-1} dz_l = -\frac{1}{T} \text{Res}_{p_{\infty-}} \frac{1}{z - z_l} z_{l+1} dz_l \end{aligned} \quad (132)$$

**Remark:** we have:

$$\mu = -T \ln \gamma \tilde{\gamma} + \sum_i \epsilon_i \eta_i + \text{Res}_{p_{\infty+}} \left( \sum_{k=0}^{\mathcal{N}} V_k(z_k) - \sum_{k=1}^{\mathcal{N}} z_{k-1} z_k \right) dS \quad (133)$$

where  $dS$  and  $\gamma$  and  $\tilde{\gamma}$  and  $\eta_i$  are defined in eq.(96) and eq.(100).

**Proof of eq.(133):** Consider the functions:

$$\phi_0(p) := \int_{p_{\infty+}}^p \left( \frac{T}{z_0} - z_{-1} \right) dz_0, \quad \phi_{\mathcal{N}}(p) := \int_{p_{\infty-}}^p \left( \frac{T}{z_{\mathcal{N}}} - z_{\mathcal{N}+1} \right) dz_{\mathcal{N}} \quad (134)$$

They satisfy:

$$\text{Res}_{p_{\infty+}} \phi_0 dS = 0, \quad \text{Res}_{p_{\infty-}} \phi_{\mathcal{N}} dS = 0 \quad (135)$$

and they have the following discontinuities along  $\mathcal{A}_i$  or  $\mathcal{B}_i$ :

$$\text{disc}_{\mathcal{A}_i} \phi_0 = -\text{disc}_{\mathcal{A}_i} \phi_{\mathcal{N}} = \Gamma_i, \quad \text{disc}_{\mathcal{B}_i} \phi_0 = -\text{disc}_{\mathcal{B}_i} \phi_{\mathcal{N}} = 2i\pi \epsilon_i \quad (136)$$

From eq.(112) and eq.(96), we have:

$$\begin{aligned} \mu &= \text{Res}_{p_{\infty+}} \mu dS = \text{Res}_{p_{\infty+}} \phi_0 dS + \text{Res}_{p_{\infty+}} \phi_{\mathcal{N}} dS - T \ln \gamma \tilde{\gamma} \\ &\quad + \text{Res}_{p_{\infty+}} \left( \sum_{k=0}^{\mathcal{N}} V_k(z_k(p)) - \sum_{k=1}^{\mathcal{N}} z_{k-1}(p) z_k(p) \right) dS \end{aligned} \quad (137)$$

We need to compute  $\text{Res}_{p_{\infty+}} \phi_{\mathcal{N}} dS$ , we use Riemann's bilinear identity:

$$\begin{aligned} \text{Res}_{p_{\infty+}} \phi_{\mathcal{N}} dS &= -\text{Res}_{p_{\infty-}} \phi_{\mathcal{N}} dS + \sum_i \frac{1}{2i\pi} \oint_{\mathcal{A}_i} \text{disc}_{\mathcal{A}_i} \phi_{\mathcal{N}} dS - \sum_i \frac{1}{2i\pi} \oint_{\mathcal{B}_i} \text{disc}_{\mathcal{B}_i} \phi_{\mathcal{N}} dS \\ &= \sum_i \eta_i \epsilon_i \end{aligned} \quad (138)$$

QED.

## 5.2 Derivatives with respect to $\epsilon_i$

We have:

$$\frac{d\mu}{d\epsilon_i} = \eta_i \quad (139)$$

indeed, from eq.(112), we have:

$$\frac{\partial \mu}{\partial \epsilon_i} = - \int_{p_{\infty+}}^p \frac{\partial z_{-1}}{\partial \epsilon_i} \Big|_{z_0} dz_0 - \int_{p_{\infty-}}^p \frac{\partial z_{\mathcal{N}+1}}{\partial \epsilon_i} \Big|_{z_{\mathcal{N}}} dz_{\mathcal{N}} \quad (140)$$

using eq.(108) we have:

$$\frac{\partial \mu}{\partial \epsilon_i} = 2i\pi \int_{p_{\infty+}}^{p_{\infty-}} du_i = \eta_i \quad (141)$$

We have:

$$\frac{dF^{(0)}}{d\epsilon_i} = \Gamma_i \quad (142)$$

Indeed, from eq.(111), and using eq.(108), we have:

$$\begin{aligned}
4 \frac{\partial F^{(0)}}{\partial \epsilon_i} &= 2T\eta_i + 2\Gamma_i + 2 \sum_j \epsilon_j \tau_{i,j} + \sum_k \operatorname{Res}_{p_{\infty+}} (V_k(z_k) - z_k z_{k+1}) du_i \\
&\quad - \sum_k \operatorname{Res}_{p_{\infty-}} (V_k(z_k) - z_k z_{k-1}) du_i
\end{aligned}
\tag{143}$$

Let us introduce the multivalued function:

$$u_i(p) := \int_{p_{\infty+}}^p du_i \tag{144}$$

its discontinuities along the  $\mathcal{A}$  and  $\mathcal{B}$  cycles are:

$$\operatorname{disc}_{\mathcal{A}_j} u_i = \tau_{i,j} , \quad \operatorname{disc}_{\mathcal{B}_j} u_i = \delta_{i,j} \tag{145}$$

After an integration by parts, we have:

$$\begin{aligned}
4 \frac{\partial F}{\partial \epsilon_i} &= 2T\eta_i + 2\Gamma_i + 2 \sum_j \epsilon_j \tau_{i,j} - \sum_k \operatorname{Res}_{p_{\infty+}} u_i (V'_k(z_k) dz_k - z_{k+1} dz_k - z_k dz_{k+1}) \\
&\quad + \sum_k \operatorname{Res}_{p_{\infty-}} u_i (V'_k(z_k) dz_k - z_{k-1} dz_k - z_k dz_{k-1}) \\
&= 2T\eta_i + 2\Gamma_i + 2 \sum_j \epsilon_j \tau_{i,j} - \sum_k \operatorname{Res}_{p_{\infty+}} u_i (z_{k-1} dz_k - z_k dz_{k+1}) \\
&\quad + \sum_k \operatorname{Res}_{p_{\infty-}} u_i (z_{k+1} dz_k - z_k dz_{k-1}) \\
&= 2T\eta_i + 2\Gamma_i + 2 \sum_j \epsilon_j \tau_{i,j} - \operatorname{Res}_{p_{\infty+}} u_i z_{-1} dz_0 + \operatorname{Res}_{p_{\infty+}} u_i z_{\mathcal{N}} dz_{\mathcal{N}+1} \\
&\quad + \operatorname{Res}_{p_{\infty-}} u_i z_{\mathcal{N}+1} dz_{\mathcal{N}} - \operatorname{Res}_{p_{\infty-}} u_i z_0 dz_{-1} \\
&= 2T\eta_i + 2\Gamma_i + 2 \sum_j \epsilon_j \tau_{i,j} + \operatorname{Res}_{p_{\infty+}} u_i z_{\mathcal{N}} dz_{\mathcal{N}+1} + T u_i(p_{\infty-}) - \operatorname{Res}_{p_{\infty-}} u_i z_0 dz_{-1}
\end{aligned}
\tag{146}$$

using Riemann's bilinear identity and eq.(133), we find eq.(142).

### 5.3 Derivatives with respect to $T$

We have:

$$\frac{\partial \mu}{\partial T} = -\ln \gamma \tilde{\gamma} \tag{147}$$

indeed, from eq.(112), we have:

$$\frac{\partial \mu}{\partial T} = \int_{p_{\infty+}}^p \left( \frac{1}{z_0} - \frac{\partial z_{-1}}{\partial T} \Big|_{z_0} \right) dz_0 + \int_{p_{\infty-}}^p \left( \frac{1}{z_{\mathcal{N}}} - \frac{\partial z_{\mathcal{N}+1}}{\partial T} \Big|_{z_{\mathcal{N}}} \right) dz_{\mathcal{N}} - \ln z_0 - \ln z_{\mathcal{N}} \quad (148)$$

using eq.(103) and eq.(99), we have:

$$\begin{aligned} \frac{\partial \mu}{\partial T} &= \int_{p_{\infty+}}^p \left( \frac{dz_0}{z_0} - \frac{d\Lambda}{\Lambda} \right) + \int_{p_{\infty-}}^p \left( \frac{dz_{\mathcal{N}}}{z_{\mathcal{N}}} + \frac{d\Lambda}{\Lambda} \right) - \ln z_0 - \ln z_{\mathcal{N}} \\ &= \ln \frac{z_0}{\Lambda} - \ln \gamma + \ln z_{\mathcal{N}} \Lambda - \ln \tilde{\gamma} - \ln z_0 - \ln z_{\mathcal{N}} = -\ln \gamma \tilde{\gamma} \end{aligned} \quad (149)$$

We have:

$$\frac{dF^{(0)}}{dT} = \mu - T \left( \frac{\mathcal{N} + 3}{2} + \mathcal{N} \ln T \right) \quad (150)$$

Indeed, from eq.(111), and using eq.(103), we have:

$$\begin{aligned} 4 \frac{\partial F}{\partial T} &= 2\mu + 2T \frac{\partial \mu}{\partial T} + 2 \sum_i \epsilon_i \eta_i + \sum_k \text{Res}_{p_{\infty+}} (V_k(z_k) - z_k z_{k+1}) dS \\ &\quad - \sum_k \text{Res}_{p_{\infty-}} (V_k(z_k) - z_k z_{k-1}) dS - 4T(1 + \mathcal{N} \ln T) - 2\mathcal{N}T \end{aligned} \quad (151)$$

using eq.(147) and eq.(133), we find eq.(150).

**Remark:** The derivative of the free energy wrt to T can be computed directly from eq.(1):

$$-N^2 \frac{\partial(F/T^2)}{\partial T} \Big|_{\epsilon_i/T} = \frac{N}{T^2} \left\langle \text{tr} \sum_{k=0}^{\mathcal{N}} V_k(M_k) - \sum_{k=1}^{\mathcal{N}} M_k M_{k-1} \right\rangle \quad (152)$$

The matrix integral is defined for fixed filling fractions, i.e. for fixed  $\epsilon_i/T$ . We can thus write:

$$T^2 \frac{\partial(F/T^2)}{\partial T} + \sum_i \frac{\epsilon_i}{T} \frac{\partial F}{\partial \epsilon_i} = -\frac{1}{N} \left\langle \text{tr} \sum_{k=0}^{\mathcal{N}} V_k(M_k) - \sum_{k=1}^{\mathcal{N}} M_k M_{k-1} \right\rangle \quad (153)$$

Using the change of variable  $\delta M_k = M_k$ , we get the loop equation:

$$2T = \frac{1}{N} \langle \text{tr} M_k V'_k(M_k) - M_k(M_{k-1} + M_{k+1}) \rangle \quad (154)$$

which implies:

$$-T^2 \frac{\partial(F/T^2)}{\partial T} - \frac{1}{T} \sum_i \epsilon_i \Gamma_i = \frac{1}{N} \left\langle \text{tr} \sum_{k=0}^{\mathcal{N}} (V_k(M_k) - \frac{1}{2} M_k V'_k(M_k)) \right\rangle + (\mathcal{N}+1)T \quad (155)$$

i.e.

$$2\frac{F}{T} - \frac{\partial F}{\partial T} - \frac{1}{T} \sum_i \epsilon_i \Gamma_i = \sum_{k=0}^{\mathcal{N}} \text{Res} (V_k(M_k) - \frac{1}{2} M_k V'_k(M_k)) W_k dz_k + (\mathcal{N}+1)T \quad (156)$$

which is equivalent to eq.(150).

Note that eq.(154), is nothing but the infinitesimal version of the rescaling  $M_k \rightarrow \alpha M_k$ . In particular one can choose  $\alpha = \sqrt{T}$ , and get directly from eq.(1):

$$F(g_{k,j}, \tilde{g}_{k,j}, \epsilon_i, T) = T^2 F(g_{k,j} T^{j/2-1}, \tilde{g}_{k,j} T^{j/2-1}, \epsilon_i/T, 1) - \frac{\mathcal{N}+1}{2} T^2 \ln T \quad (157)$$

Taking the derivative of eq.(157) with respect to  $T$ , gives again eq.(150).

## 6 Other observables, leading order

### 6.1 Resolvents, leading order

So far, we have found that  $W_0(z_0)$  and  $W_{\mathcal{N}}(z_{\mathcal{N}})$  obey algebraic equations. We have been able to determine the resolvents for the matrices at the extremities of the chain, but not for intermediate matrices. We are going to determine  $W_k(z_k)$  for  $0 \leq k \leq \mathcal{N}$ .

We start from eq.(132):

$$W_k(z) = -\frac{1}{2i\pi T} \oint_{\mathcal{C}_k} \frac{z_{k-1}}{z - z_k} dz_k \quad (158)$$

where the integration contour is  $\mathcal{C}_k$  defined in 4.3. Indeed, the residue is a contour integral around  $p_{\infty+}$ , and since equation eq.(132) was derived formally for large  $z$  (i.e. order by order in the large  $z$  expansion), we have to assume that the contour of integration encloses all the  $r_k$  solutions of

$z_k(p) = z$ . Moreover, the residue is the sum of poles at the  $p_{+j,k}(z_k)$ 's and at  $p_{\infty+}$ :

$$\begin{aligned}
\frac{1}{2i\pi} \oint_{\mathcal{C}_k} \frac{z_{k-1}}{z_k - z} dz_k &= \sum_{j=1}^{r_k} z_{k-1}(p_{+j,k}(z_k)) + \frac{1}{2i\pi} \oint_{\mathcal{C}_k} z_{k-1} \frac{dz_k}{z_k} \\
&= \sum_{j=1}^{r_k} z_{k-1}(p_{+j,k}(z_k)) + \frac{1}{2i\pi} \oint_{\mathcal{C}_k} z_{k-1} \frac{V''(z_{k-1}) dz_{k-1} - dz_{k-2}}{V'_{k-1}(z_{k-1}) - z_{k-2}} \\
&= \sum_{j=1}^{r_k} z_{k-1}(p_{+j,k}(z_k)) + \frac{1}{2i\pi} \oint_{\mathcal{C}_k} \frac{V''(z_{k-1}) z_{k-1}}{V'_{k-1}(z_{k-1})} dz_{k-1} \\
&= \sum_{j=1}^{r_k} z_{k-1}(p_{+j,k}(z_k)) + \frac{g_{k-1,d_{k-1}}}{g_{k-1,d_{k-1}+1}} \tag{159}
\end{aligned}$$

Therefore:

$$TW_k(z) = \frac{g_{k-1,d_{k-1}}}{g_{k-1,d_{k-1}+1}} + \sum_{j=1}^{r_k} z_{k-1}(p_{+j,k}(z)) = \frac{g_{k+1,d_{k+1}}}{g_{k+1,d_{k+1}+1}} + \sum_{j=1}^{s_k} z_{k+1}(p_{-j,k}(z)) \tag{160}$$

Notice that we have:

$$\sum_{j=-s_k}^{r_k} z_{k-1}(p_{j,k}(z)) = s_k V'_k(z) + \frac{g_{k+1,d_{k+1}}}{g_{k+1,d_{k+1}+1}} - \frac{g_{k-1,d_{k-1}}}{g_{k-1,d_{k-1}+1}} \tag{161}$$

## 6.2 2-loop functions, leading order

From eq.(87) and eq.(160) We find ( $l \leq k$ ):

$$\boxed{\left\langle \text{Tr} \frac{1}{z_k - M_k} \text{Tr} \frac{1}{z_l - M_l} \right\rangle_c dz_k dz_l = - \sum_{i=1}^{r_k} \sum_{j=1}^{r_l} B(p_{+i,k}(z_k), p_{+j,l}(z_l))} \tag{162}$$

## 6.3 2-point one loop functions

### 6.3.1 $W_{k,k+1}$

Define the following polynomial of two variables:

$$Q_{k,k+1}(z_k, z_{k+1}) = \text{Res}_{\infty} dz_0 \dots dz_{k-1} dz_{k+2} \dots dz_N \frac{E(z_0, \dots, z_k, z_{k+1}, \dots, z_N)}{\prod_{j=0}^N (V_j(z_j) - z_{j+1} - z_{j-1})}$$

$$\begin{aligned}
&= E(\underline{Z}_0, \dots, \underline{Z}_{k-1}, z_k, z_{k+1}, \underline{Z}_{k+2}, \dots, \underline{Z}_{\mathcal{N}}) \\
&= (-1)^{r_k} \prod_{j=0}^{k-1} g_{j,d_j+1}^{r_j} \prod_{j=k+1}^{\mathcal{N}} g_{j,d_j+1}^{s_j} \prod_{0 \neq j = -s_k}^{r_k} (z_{k+1} - z_{k+1}(p_{j,k}(z_k))) \\
&= (-1)^{s_{k+1}} \prod_{j=0}^k g_{j,d_j+1}^{r_j} \prod_{j=k+2}^{\mathcal{N}} g_{j,d_j+1}^{s_j} \prod_{0 \neq j = -s_{k+1}}^{r_{k+1}} (z_k - z_k(p_{j,k+1}(z_{k+1})))
\end{aligned} \tag{163}$$

where the  $\underline{Z}_j(z_k, z_{k+1})$  are defined by:

$$\begin{cases} \underline{Z}_k(z_k, z_{k+1}) = z_k, & \underline{Z}_{k+1}(z_k, z_{k+1}) = z_{k+1} \\ \underline{Z}_{j+1}(z_k, z_{k+1}) = V'_l(\underline{Z}_j(z_k, z_{k+1})) - \underline{Z}_{j-1}(z_k, z_{k+1}) & j > k+1 \\ \underline{Z}_{j-1}(z_k, z_{k+1}) = V'_l(\underline{Z}_j(z_k, z_{k+1})) - \underline{Z}_{j+1}(z_k, z_{k+1}) & j < k \end{cases} \tag{164}$$

We conjecture (proved in appendix B for  $k = 0$ ):

$$\begin{aligned}
1 - W_{k,k+1}(z_k, z_{k+1}) &= \frac{Q_{k,k+1}(z_k, z_{k+1}) / \prod_{j=0}^{k-1} g_{j,d_j+1}^{r_j} \prod_{j=k+2}^{\mathcal{N}} g_{j,d_j+1}^{s_j}}{\prod_{j=1}^{r_k} (z_{k+1}(p_{+j,k}(z_k)) - z_{k+1}) \prod_{j=1}^{s_{k+1}} (z_k(p_{-j,k+1}(z_{k+1})) - z_k)} \\
&= g_{k,d_k+1}^{r_k} \frac{\prod_{j=1}^{r_{k+1}} (z_k - z_k(p_{+j,k+1}(z_{k+1})))}{\prod_{j=1}^{r_k} (z_{k+1}(p_{+j,k}(z_k)) - z_{k+1})} \\
&= g_{k+1,d_{k+1}+1}^{s_{k+1}} \frac{\prod_{j=1}^{s_k} (z_{k+1} - z_{k+1}(p_{-j,k}(z_k)))}{\prod_{j=1}^{s_{k+1}} (z_k(p_{-j,k+1}(z_{k+1})) - z_k)}
\end{aligned} \tag{165}$$

**Remark:**

$$Q_{k,k+1}(z_k, z_{k+1}) = Q_{k-1,k}(V'_k(z_k) - z_{k+1}, z_k) \tag{166}$$

and we may conjecture that the spectral curves of the differential systems defined in the appendix of [5] are:

$$\det(z_{k+1} \mathbf{1}_{r_k+s_k} - \mathcal{D}_k(z_k)) = (-1)^{r_k} \frac{Q_{k,k+1}(z_k, z_{k+1})}{\prod_{j=0}^{k-1} g_{j,d_j+1}^{r_j} \prod_{j=k+1}^{\mathcal{N}} g_{j,d_j+1}^{s_j}} \tag{167}$$

so that the property eq.(166) would be nothing but the duality discovered in [5].

### 6.3.2 The function $U$ in the large $N$ limit

Since the function  $U$  appears in the RHS of the master loop equation, it is important to be able to compute it in the large  $N$  limit.

Define (notice that  $U_1 = U$ ):

$$U_k(z_0, z_k, \dots, z_N) := \text{Pol}_{z_1, \dots, z_N} W(z_0, \dots, z_N) f_{k,N}(z_k, \dots, z_N) \prod_{j=1}^{k-1} z_j \quad (168)$$

We shall prove that:

$$\begin{aligned} U_k(z_0, z_k, \dots, z_N) &= H_{k,N}(z_k, \dots, z_N) W_0(z_0) \\ &\quad - \sum_{l=k}^N \frac{P(Z_0, \dots, Z_{l-1}, z_l, \dots, z_N) - P(Z_0, \dots, Z_l, z_{l+1}, \dots, z_N)}{z_l - Z_l} H_{k,l-1}(z_k, \dots, z_{l-1}) \\ &= V'_N(z_N) - z_{N-1} \\ &\quad + \sum_{l=0}^{N-1} \frac{E(Z_0, \dots, Z_{l-1}, z_l, \dots, z_N) - E(Z_0, \dots, Z_l, z_{l+1}, \dots, z_N)}{z_l - Z_l} H_{k,l-1}(z_k, \dots, z_{l-1}) \end{aligned} \quad (169)$$

where

$$H_{k,l}(z_k, \dots, z_l) := \text{Pol}_{z_0, \dots, z_N} \prod_{j=0}^N \frac{1}{z_j - Z_j} f_{k,l}(z_k, \dots, z_l) \prod_{j < k} z_j \prod_{j > l} z_j \quad (170)$$

i.e.  $H_{k,l}$  is a polynomial in  $z_k, \dots, z_l$ , and satisfies:

$$H_{k+1,k} = 1, \quad H_{k,l} = 0 \text{ if } k > l+1, \quad H_{k,l} = \frac{V'_k(z_k) - V'_k(Z_k)}{z_k - Z_k} H_{k+1,l} - H_{k+2,l} \quad (171)$$

In particular for  $k=1$  eq.(169) reduces to:

$$\begin{aligned} U(z_0, \dots, z_N) &= V'_N(z_N) - z_{N-1} \\ &\quad + \sum_{k=0}^{N-1} \frac{E(Z_0, \dots, Z_k, z_{k+1}, \dots, z_N) - E(Z_0, \dots, Z_{k+1}, z_{k+2}, \dots, z_N)}{z_{k+1} - Z_{k+1}} H_{1,k}(z_1, \dots, z_k) \\ &= H_{1,N}(z_1, \dots, z_N) W_0(z_0) \\ &\quad - \sum_{k=0}^{N-1} \frac{P(Z_0, \dots, Z_k, z_{k+1}, \dots, z_N) - P(Z_0, \dots, Z_{k+1}, z_{k+2}, \dots, z_N)}{z_{k+1} - Z_{k+1}} H_{1,k}(z_1, \dots, z_k) \end{aligned} \quad (172)$$

**Proof of eq.(169):**

Let us define:

$$A_k(z_0, z_k, \dots, z_N) := \text{Pol}_{z_1, \dots, z_{k-1}, z_{k+1}, \dots, z_N} \prod_{j=1}^{k-1} z_j f_{k+1,N}(z_{k+1}, \dots, z_N) W(z_0, \dots, z_N) \quad (173)$$



$$S_k(z_0, z_k, \dots, z_{\mathcal{N}}) := \frac{\text{Pol}_{z_0, \dots, z_{k-1}, z_{k+1}, \dots, z_{\mathcal{N}}} f_{0,k-1}(z_0, \dots, z_{k-1})}{f_{k+1, \mathcal{N}}(z_{k+1}, \dots, z_{\mathcal{N}})} W(z_0, \dots, z_{\mathcal{N}}) \quad (174)$$

$$T_k(z_0, z_k, \dots, z_{\mathcal{N}}) := \frac{\text{Pol}_{z_0, \dots, z_{\mathcal{N}}} f_{0,k-1}(z_0, \dots, z_{k-1}) V'_k(z_k)}{f_{k+1, \mathcal{N}}(z_{k+1}, \dots, z_{\mathcal{N}})} W(z_0, \dots, z_{\mathcal{N}}) \quad (175)$$

From eq.(13) we have:

$$U_k = \text{Pol}_{z_k} V'_k A_k - U_{k+2} \quad (176)$$

and from eq.(29), we have:

$$(z_k - Z_k) A_k = U_{k+1} - S_k(Z_0, \dots, Z_{k-1}, z_k, \dots, z_{\mathcal{N}}) \quad (177)$$

That implies:

$$U_k = \frac{V'_k(z_k) - V'_k(Z_k)}{z_k - Z_k} U_{k+1} - U_{k+2} - \frac{T_k(Z_0, \dots, Z_{k-1}, z_k, \dots, z_{\mathcal{N}}) - T_k(Z_0, \dots, Z_k, z_{k+1}, \dots, z_{\mathcal{N}})}{z_k - Z_k} \quad (178)$$

It is easy to prove (by recursion), that  $f_{0, \mathcal{N}} - f_{0,k-1} V'_k(z_k) f_{k+1, \mathcal{N}}$  is linear in  $z_k$ , in other words we can write:

$$f_{0, \mathcal{N}} = f_{0,k-1} V'_k(z_k) f_{k+1, \mathcal{N}} + z_k B_k \quad (179)$$

where  $B_k$  is independent of  $z_k$ , therefore:

$$T_k(z_0, \dots, z_k, \dots, z_{\mathcal{N}}) - P(z_0, \dots, z_k, \dots, z_{\mathcal{N}}) = \text{independent of } z_k \quad (180)$$

and thus eq.(178) can be rewritten:

$$U_k = \frac{V'_k(z_k) - V'_k(Z_k)}{z_k - Z_k} U_{k+1} - U_{k+2} - \frac{P(Z_0, \dots, Z_{k-1}, z_k, \dots, z_{\mathcal{N}}) - P(Z_0, \dots, Z_k, z_{k+1}, \dots, z_{\mathcal{N}})}{z_k - Z_k} \quad (181)$$

From which, the initial conditions  $U_{\mathcal{N}+1} = W_0$ ,  $U_{\mathcal{N}+2} = 0$  (easily derived) imply eq.(169).

### 6.3.3 The extremities correlator $W_{0,\mathcal{N}}$

From eq.(172) and eq.(41), we derive:

$$W_{0,\mathcal{N}}(z_0, z_{\mathcal{N}}) = \sum_{k=0}^{\mathcal{N}-1} \frac{E(Z_0, \dots, Z_k, \tilde{Z}_{k+1}, \dots, \tilde{Z}_{\mathcal{N}})}{(Z_k - \tilde{Z}_k)(Z_{k+1} - \tilde{Z}_{k+1})} \quad (182)$$

where:

$$Z_k := z_k(p_{+1,0}(z_0)) , \quad \tilde{Z}_k := z_k(p_{-1,\mathcal{N}}(z_{\mathcal{N}})) \quad (183)$$

## 7 Subleading expansion

The aim of this section is to generalize the calculation of [10, 11], and compute the next to leading  $1/N^2$  term in the topological expansion. In this purpose, we expand the various observables in a  $1/N^2$  power series:

$$Z_k(z_0) = Z_k^{(0)}(z_0) + \frac{1}{N^2} Z_k^{(1)}(z_0) + \dots \quad (184)$$

$$P(z_0, \dots, z_{\mathcal{N}}) = P^{(0)}(z_0, \dots, z_{\mathcal{N}}) + \frac{1}{N^2} P^{(1)}(z_0, \dots, z_{\mathcal{N}}) + \dots \quad (185)$$

and so on. Then we expand eq.(28) to order  $1/N^2$ :

$$\sum_{k=1}^{\mathcal{N}} Z_k^{(1)}(z_0) E_k - P^{(1)}(z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)}) = U_{;0}^{(0)}(z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)}; z_0) \quad (186)$$

where:

$$E_k := \left. \frac{\partial E^{(0)}(z_0, \dots, z_{\mathcal{N}})}{\partial z_k} \right|_{z_j = Z_j^{(0)}} \quad (187)$$

From the definition of the  $Z_k$ 's eq.(11), we easily find:

$$Z_k^{(1)}(z_0) = \overline{H}_{1,k-1} Z_1^{(1)}(z_0) \quad (188)$$

where (the  $H_{k,l}$ 's have been defined in eq.(170)):

$$\overline{H}_{1,k} := H_{1,k}(Z_1^{(0)}, \dots, Z_k^{(0)}) \quad (189)$$

We have:

$$\overline{H}_{1,-1} = 0 , \quad \overline{H}_{1,0} = 1 , \quad \overline{H}_{1,k} = V_k''(Z_k^{(0)}) \overline{H}_{1,k-1} - \overline{H}_{1,k-2} \quad (190)$$

i.e.

$$\overline{H}_{1,k} = \frac{\partial^k f_{1,k}(z_1, \dots, z_k)}{\partial z_1 \dots \partial z_k} \Big|_{z_j = Z_j^{(0)}} \quad (191)$$

Therefore eq.(185) reads:

$$Z_1^{(1)}(z_0) = \frac{P^{(1)}(z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)}) + U_{;0}^{(0)}(z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)}; z_0)}{\sum_{k=1}^{\mathcal{N}} \overline{H}_{1,k-1} E_k} \quad (192)$$

Similarly to what was done in [10, 11], we determine the polynomial  $P^{(1)}$ , by the condition that  $Z_1^{(1)}(z_0)$  has singularities only at the endpoints  $e_{0,i}$  in the  $z_0$ -sheet. That condition is sufficient to determine all the unknown coefficients of  $P^{(1)}$ . Now, let us compute  $U_{;0}^{(0)}$ .

## 7.1 Computation of $U_{;0}(z_0, \dots, z_{\mathcal{N}}; x)$

Notice that there is no explicit dependence on  $V_0$  in eq.(26), therefore:

$$U_{;0}(z_0, \dots, z_{\mathcal{N}}; z) = \frac{\partial U(z_0, \dots, z_{\mathcal{N}})}{\partial V_0(z)} \quad (193)$$

For  $k \geq 1$ , write  $z_k = Z_k^{(0)} + \zeta_k$ , and Taylor expand eq.(172) in the  $\zeta$ 's:

$$\begin{aligned} U(z_0, \dots, z_{\mathcal{N}}) &= V'_{\mathcal{N}}(z_{\mathcal{N}}) - z_{\mathcal{N}-1} \\ &\quad - \sum_{k=1}^{\mathcal{N}} \left( E_k + \frac{1}{2} \zeta_k E_{k,k} + \sum_{i=k+1}^{\mathcal{N}} \zeta_i E_{k,i} \right) \\ &\quad \left( \overline{H}_{1,k-1} + \sum_{i=1}^{k-1} \zeta_i \overline{H}_{1,k-1;i} \right) \end{aligned} \quad (194)$$

where  $E_k$  was defined in eq.(187), and  $E_{k,l}$  is:

$$E_{k,l} := \frac{\partial^2 E^{(0)}(z_0, \dots, z_{\mathcal{N}})}{\partial z_k \partial z_l} \Big|_{z_j = Z_j^{(0)}} \quad (195)$$

$\overline{H}_{1,k}$  was defined in eq.(189) and:

$$\overline{H}_{1,k;i} := \frac{\partial H_{1,k}(z_1, \dots, z_k)}{\partial z_i} \Big|_{z_j = Z_j^{(0)}} \quad (196)$$

From eq.(170) and eq.(189) we derive:

$$\overline{H}_{1,k;i} = \frac{1}{2} V_i'''(Z_i^{(0)}) \overline{H}_{1,i-1} \overline{H}_{i+1,k} \quad (197)$$

Then take the  $\partial/\partial V_0(z)$  derivative, using  $\partial\zeta_k/\partial V_0 = -\partial Z_k/\partial V_0$ , and take the  $\zeta \rightarrow 0$  limit:

$$\begin{aligned} U_{;0}^{(0)}(Z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)}; z) &= \sum_{k=1}^{\mathcal{N}} \left( \frac{\partial E_k}{\partial V_0(z)} - \frac{\partial Z_k^{(0)}}{\partial V_0(z)} \frac{E_{k,k}}{2} - \sum_{j>k} \frac{\partial Z_j^{(0)}}{\partial V_0(z)} E_{k,j} \right) \overline{H}_{1,k-1} \\ &+ \sum_{k=1}^{\mathcal{N}} E_k \left( \frac{\partial \overline{H}_{1,k-1}}{\partial V_0(z)} - \sum_{j=1}^{k-1} \frac{\partial Z_j^{(0)}}{\partial V_0(z)} \overline{H}_{1,k-1;j} \right) \end{aligned} \quad (198)$$

From eq.(11), one easily derives:

$$\frac{\partial Z_{k+1}^{(0)}}{\partial V_0(z)} = \frac{\partial Z_k^{(0)}}{\partial V_0(z)} V_k''(Z_k^{(0)}) - \frac{\partial Z_{k-1}^{(0)}}{\partial V_0(z)} = \overline{H}_{1,k} \frac{\partial Z_1^{(0)}}{\partial V_0(z)} \quad (199)$$

and from eq.(189)

$$\frac{\partial \overline{H}_{1,k-1}}{\partial V_0(z)} = \sum_{j=1}^{k-1} \frac{\partial Z_j^{(0)}}{\partial V_0(z)} V_j'''(Z_j^{(0)}) \overline{H}_{1,j-1} \overline{H}_{j+1,k-1} = 2 \sum_{j=1}^{k-1} \frac{\partial Z_j^{(0)}}{\partial V_0(z)} \overline{H}_{1,k-1;j} \quad (200)$$

## 7.2 Computation of $\partial E_k/\partial V_0(z)$

Let us define  $\overline{Z}_j(z_0, z_1)$ :

$$\overline{Z}_0(z_0, z_1) := z_0, \quad \overline{Z}_1(z_0, z_1) := z_1, \quad \overline{Z}_{j+1}(z_0, z_1) = V_j'(\overline{Z}_j(z_0, z_1)) - \overline{Z}_{j-1}(z_0, z_1) \quad (201)$$

We have:

$$\left. \frac{\partial \overline{Z}_j(z_0, z_1)}{\partial z_1} \right|_{z_1=Z_1^{(0)}} = \overline{H}_{1,j-1}, \quad \left. \frac{\partial^2 \overline{Z}_k(z_0, z_1)}{\partial z_1^2} \right|_{z_1=Z_1^{(0)}} = \sum_{j=1}^{k-1} V_j'''(Z_j^{(0)}) \overline{H}_{1,j-1}^2 \overline{H}_{j+1,k-1} \quad (202)$$

Consider the polynomials defined in eq.(163):

$$Q_{0,1}(z_0, z_1) := E^{(0)}(\overline{Z}_0(z_0, z_1), \dots, \overline{Z}_{\mathcal{N}}(z_0, z_1))$$

$$= C(z_1 - Z_1^{(0)}(z_0)) \prod_{j=1}^{s_0} (z_1 - z_1(p_{-j,0}(z_0))) \quad (203)$$

(where  $C = -\prod_{j=1}^{\mathcal{N}} g_{j,d_j+1}^{s_j}$ ), and take its derivative (using eq.(202)):

$$\left. \frac{\partial Q_{0,1}(z_0, z_1)}{\partial z_1} \right|_{z_1=Z_1^{(0)}} = \sum_{j=1}^{\mathcal{N}} E_j \bar{H}_{1,j-1} = C \prod_{j=1}^{s_0} (Z_1^{(0)} - z_1(p_{-j,0}(z_0))) \quad (204)$$

That implies:

$$\begin{aligned} & \sum_{j=1}^{\mathcal{N}} \frac{\partial E_j}{\partial V_0(z)} \bar{H}_{1,j-1} + \sum_{j=1}^{\mathcal{N}} E_j \frac{\partial \bar{H}_{1,j-1}}{\partial V_0(z)} \\ &= \sum_{l=1}^{s_0} \left( \frac{\partial Z_1^{(0)}}{\partial V_0(z)} - \frac{\partial z_1(p_{-l,0}(z_0))}{\partial V_0(z)} \right) \frac{\sum_{j=1}^{\mathcal{N}} E_j \bar{H}_{1,j-1}}{Z_1^{(0)} - z_1(p_{-l,0}(z_0))} \end{aligned} \quad (205)$$

Plugging that into eq.(198), we get:

$$\begin{aligned} U_{;0}^{(0)}(Z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)}; z) &= \sum_{l=1}^{s_0} \left( \frac{\partial Z_1^{(0)}}{\partial V_0(z)} - \frac{\partial z_1(p_{-l,0}(z_0))}{\partial V_0(z)} \right) \frac{\sum_{j=1}^{\mathcal{N}} E_j \bar{H}_{1,j-1}}{Z_1^{(0)} - z_1(p_{-l,0}(z_0))} \\ &\quad - \sum_{k=1}^{\mathcal{N}} \frac{\partial Z_k^{(0)}}{\partial V_0(z)} \frac{E_{k,k}}{2} \bar{H}_{1,k-1} - \sum_{k=1}^{\mathcal{N}} \sum_{j>k} \frac{\partial Z_j^{(0)}}{\partial V_0(z)} E_{k,j} \bar{H}_{1,k-1} \\ &\quad - \sum_{k=1}^{\mathcal{N}} \sum_{j=1}^{k-1} \frac{\partial Z_j^{(0)}}{\partial V_0(z)} E_k \bar{H}_{1,k-1;j} \\ &= \sum_{l=1}^{s_0} \left( \frac{\partial Z_1^{(0)}}{\partial V_0(z)} - \frac{\partial z_1(p_{-l,0}(z_0))}{\partial V_0(z)} \right) \frac{\sum_{j=1}^{\mathcal{N}} E_j \bar{H}_{1,j-1}}{Z_1^{(0)} - z_1(p_{-l,0}(z_0))} \\ &\quad - \frac{\partial Z_1^{(0)}}{\partial V_0(z)} \sum_{k=1}^{\mathcal{N}} \left( \frac{E_{k,k}}{2} \bar{H}_{1,k-1}^2 + \sum_{j>k} E_{k,j} \bar{H}_{1,k-1} \bar{H}_{1,j-1} \right) \\ &\quad - \frac{1}{2} \frac{\partial Z_1^{(0)}}{\partial V_0(z)} \sum_{k=1}^{\mathcal{N}} \sum_{j=1}^{k-1} E_k V_j'''(Z_j^{(0)}) \bar{H}_{1,j-1}^2 \bar{H}_{j+1,k-1} \\ &= - \sum_{l=1}^{s_0} \frac{\partial z_1(p_{-l,0}(z_0))}{\partial V_0(z)} \bigg|_{z_0} \frac{\left( \sum_{j=1}^{\mathcal{N}} E_j \bar{H}_{1,j-1} \right)}{Z_1^{(0)} - z_1(p_{-l,0}(z_0))} \end{aligned}$$

$$\begin{aligned}
& + \frac{\partial Z_1^{(0)}}{\partial V_0(z)} \sum_{l=1}^{s_0} \frac{\left( \sum_{j=1}^{\mathcal{N}} E_j \bar{H}_{1,j-1} \right)}{Z_1^{(0)} - z_1(p_{-l,0}(z_0))} \\
& - \frac{1}{2} \frac{\partial Z_1^{(0)}}{\partial V_0(z)} \sum_{k=1}^{\mathcal{N}} \sum_{j=1}^{\mathcal{N}} E_{k,j} \bar{H}_{1,k-1} \bar{H}_{1,j-1} \\
& - \frac{1}{2} \frac{\partial Z_1^{(0)}}{\partial V_0(z)} \sum_{k=1}^{\mathcal{N}} \sum_{j=1}^{k-1} E_k V_j'''(Z_j^{(0)}) \bar{H}_{1,j-1}^2 \bar{H}_{j+1,k-1} \quad (206)
\end{aligned}$$

Take the second derivative of eq.(203):

$$\begin{aligned}
\left. \frac{\partial^2 Q_{0,1}(z_0, z_1)}{\partial z_1^2} \right|_{z_1=Z_1^{(0)}} &= 2 \left( \sum_{j=1}^{\mathcal{N}} E_j \bar{H}_{1,j-1} \right) \sum_{l=1}^{s_0} \frac{1}{Z_1^{(0)} - z_1(p_{-l,0}(z_0))} \\
&= \sum_{k=1}^{\mathcal{N}} \sum_{j=1}^{\mathcal{N}} E_{k,j} \bar{H}_{1,j-1} \bar{H}_{1,j-1} + \sum_{k=1}^{\mathcal{N}} E_k \left. \frac{\partial^2 \bar{Z}_k(z_0, z_1)}{\partial z_1^2} \right|_{z_1=Z_1} \quad (207)
\end{aligned}$$

plugging eq.(207) and eq.(202) into eq.(206), we get:

$$\boxed{\frac{U_{;0}^{(0)}(Z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)}; z)}{\sum_{j=1}^{\mathcal{N}} E_j \bar{H}_{1,j-1}} = - \sum_{l=1}^{s_0} \frac{1}{Z_1^{(0)} - z_1(p_{-l,0}(z_0))} \left. \frac{\partial z_1(p_{-j,0}(z_0))}{\partial V_0(z)} \right|_{z_0}} \quad (208)$$

which can be compared to [11].

### 7.3 Next to leading order

Using eq.(208) into eq.(192), and using eq.(87), we get the  $1/N^2$  correction for the resolvent:

$$\begin{aligned}
Z_1^{(1)}(z_0) &= \frac{P^{(1)}(z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)})}{\sum_{k=1}^{\mathcal{N}} \bar{H}_{1,k-1} E_k} - \sum_{l=1}^{s_0} \frac{1}{Z_1^{(0)} - z_1(p_{-l,0}(z_0))} \left. \frac{\partial z_1(p_{-j,0}(z_0))}{\partial V_0(z_0)} \right|_{z_0} \\
&= \sum_{l=1}^{s_0} \frac{1}{z_1(p_{1,0}(z_0)) - z_1(p_{-l,0}(z_0))} \frac{B(p_{-l,0}(z_0), p_{1,0}(z_0))}{dz_0(p_{1,0}(z_0)) dz_0(p_{-l,0}(z_0))} \\
&\quad + \frac{P^{(1)}(z_0, Z_1^{(0)}, \dots, Z_{\mathcal{N}}^{(0)})}{\sum_{k=1}^{\mathcal{N}} \bar{H}_{1,k-1} E_k} \quad (209)
\end{aligned}$$

and we remind that  $P^{(1)}$  is completely determined by the condition that  $Z_1^{(1)}(z_0)$  has singularities only at the endpoints  $e_{0,i}$  in the  $z_0$ -sheet, and has vanishing  $\mathcal{B}$ -cycle integrals.

From there, it should be possible to extend the calculation of [10, 11] to the chain of matrices, and compute the  $1/N^2$  term in the free energy for the chain of matrices. This will be left for a later work.

## 8 Examples

### 8.1 Example: one-cut assumption (genus zero)

Let us assume that the genus is  $g = 0$ .

It was already discussed in [10] that for multimatrix models, the so called one-cut assumption should be replaced by a genus zero assumption. Indeed, in that case, the number of  $z_0$ -endpoints, i.e. the number of zeroes of  $dz_0$  is equal to the number of sheets, therefore there is exactly one cut in the physical sheet.  $\mathcal{E}$  is in one to one correspondance with the complex plane  $\mathbb{C}$ , and it is well known [16, 17] that the parametrization eq.(81) is rational, and with  $s = \frac{p-p_{\infty-}}{p-p_{\infty+}}$ , it can be written:

$$z_k(s) = \sum_{i=-s_k}^{r_k} \alpha_{k,i} s^i. \quad (210)$$

$z_k(s)$  has two poles: one pole of degree  $r_k$  at  $s = \infty$ , and one pole of degree  $s_k$  at  $s = 0$ . This parametrization is identical to the one found in [13] by the biorthogonal polynomial's method.

The set of equations:

$$\left\{ \begin{array}{l} V'_k(z_k(s)) = z_{k+1}(s) + z_{k-1}(s) \\ V'_0(z_0(s)) - z_1(s) \underset{s \rightarrow \infty}{\sim} \frac{T}{\alpha_{0,1}s} \\ V'_{\mathcal{N}}(z_{\mathcal{N}}(s)) - z_{\mathcal{N}-1}(s) \underset{s \rightarrow 0}{\sim} \frac{Ts}{\alpha_{\mathcal{N},-1}} \\ \alpha_{0,1} = \alpha_{\mathcal{N},-1} = \gamma = \tilde{\gamma} \end{array} \right. \quad (211)$$

is sufficient to determine all the  $\alpha_{k,i}$  (one is free to impose  $\alpha_{0,1} = \alpha_{\mathcal{N},-1}$  because  $s$  can be changed into any constant times  $s$ ).

The Abelian differential of the third kind is:

$$dS = -\frac{ds}{s} \quad (212)$$

The Bergmann kernel is:

$$B(s, s') = \frac{ds ds'}{(s - s')^2} \quad (213)$$

## 8.2 Example: gaussian case

Consider all potentials quadratic (i.e.  $d_k = 1$ ):

$$V_k(z_k) := \frac{g_k}{2} z_k^2 \quad (214)$$

A direct computation of the matrix integral eq.(2) gives the free energy:

$$F = F^{(0)} = \frac{T^2}{2} \ln D_{0,\mathcal{N}} - \frac{\mathcal{N} + 1}{2} T^2 \ln T \quad (215)$$

Define:

$$D_{k,l} := \det \begin{pmatrix} g_k & 1 & & & \\ 1 & \ddots & \ddots & & \\ & \ddots & \ddots & 1 & \\ & & & 1 & g_l \end{pmatrix}, \quad D_{k,k-1} := 1, \quad D_{k,k-2} := 0 \quad (216)$$

The function  $f_{k,l}$  of eq.(13) is:

$$f_{k,l} = z_k \dots z_l D_{k,l} \quad (217)$$

thus the polynomial  $P$  is a constant:

$$P(z_0, \dots, z_{\mathcal{N}}) = D_{0,\mathcal{N}} \quad (218)$$

and the polynomial  $E$  is:

$$E(z_0, \dots, z_{\mathcal{N}}) = (g_0 z_0 - z_1)(g_{\mathcal{N}} z_{\mathcal{N}} - z_{\mathcal{N}-1}) - TP \quad (219)$$

The leading order loop equations are thus:

$$z_{k+1} + z_{k-1} = g_k z_k, \quad z_{-1} z_{\mathcal{N}+1} = TP \quad (220)$$



Note that eq.(62) implies that the genus is necessarily  $g = 0$ , and thus there is a rational parametrization of the form eq.(210), with  $r_k = s_k = 1$ , namely:

$$z_k = \sqrt{\frac{T}{D_{0,\mathcal{N}}}} (D_{0,k-1}s + D_{k+1,\mathcal{N}}s^{-1}) \quad (221)$$

We find:

$$\gamma = \tilde{\gamma} = \sqrt{\frac{T}{D_{0,\mathcal{N}}}} \quad (222)$$

and:

$$\mu = T + T \ln D_{0,\mathcal{N}} - T \ln T = T - T \ln \gamma \tilde{\gamma} \quad (223)$$

and it is easy to check that eq.(111) coincides with eq.(215).

### 8.3 Loop equations of the 2 matrix model

Let us take  $T = 1$ . The 2-matrix model corresponds to  $\mathcal{N} = 1$  we have:

$$f_{0,1}(z_0, z_1) = V'_0(z_0)V'_1(z_1) - z_0z_1, \quad f_{1,1}(z_1) = V'_1(z_1) \quad (224)$$

i.e.

$$P(z_0, z_1) = \frac{1}{N} \left\langle \text{tr} \frac{V'_0(z_0) - V'_0(M_0)}{z_0 - M_0} \frac{V'_1(z_1) - V'_1(M_1)}{z_1 - M_1} \right\rangle - 1 \quad (225)$$

$$U(z_0, z_1; z) = \left\langle \text{tr} \frac{1}{z_0 - M_0} \frac{V'_1(z_1) - V'_1(M_1)}{z_1 - M_1} \text{tr} \frac{1}{z - M_0} \right\rangle_c \quad (226)$$

and the master loop equation reads:

$$(V'_0(Z_0) - Z_1)(V'_1(Z_1) - Z_0) - P(Z_0, Z_1) = \frac{1}{N^2} U(Z_0, Z_1; Z_0) \quad (227)$$

with  $Z_1 = V'_0(z_0) - W_0(z_0)$ . We recover the equation of [12, 13, 10].

### 8.4 The 3 matrix model

Let us take  $T = 1$ . The 3-matrix model corresponds to  $\mathcal{N} = 2$ , i.e.:

$$f_{0,2}(z_0, z_1, z_2) = V'_0(z_0)V'_1(z_1)V'_2(z_2) - z_0z_1V'_2(z_2) - z_2z_1V'_0(z_0) \quad (228)$$

$$f_{1,2}(z_1, z_2) = V_1'(z_1)V_2'(z_2) - z_1 z_2 \quad (229)$$

i.e.

$$\begin{aligned} P(z_0, z_1, z_2) &= \frac{1}{2N} \left\langle \text{tr} \frac{V_0'(z_0) - V_0'(M_0)}{z_0 - M_0} \frac{V_1'(z_1) - V_1'(M_1)}{z_1 - M_1} \frac{V_2'(z_2) - V_2'(M_2)}{z_2 - M_2} \right\rangle \\ &+ \frac{1}{2N} \left\langle \text{tr} \frac{V_2'(z_2) - V_2'(M_2)}{z_2 - M_2} \frac{V_1'(z_1) - V_1'(M_1)}{z_1 - M_1} \frac{V_0'(z_0) - V_0'(M_0)}{z_0 - M_0} \right\rangle \\ &- \frac{1}{N} \left\langle \text{tr} \frac{V_0'(z_0) - V_0'(M_0)}{z_0 - M_0} \right\rangle - \frac{1}{N} \left\langle \text{tr} \frac{V_2'(z_2) - V_2'(M_2)}{z_2 - M_2} \right\rangle \end{aligned} \quad (230)$$

$$\begin{aligned} U(z_0, z_1, z_2; z) &= \frac{1}{2} \left\langle \text{tr} \frac{1}{z_0 - M_0} \frac{V_1'(z_1) - V_1'(M_1)}{z_1 - M_1} \frac{V_2'(z_2) - V_2'(M_2)}{z_2 - M_2} \text{tr} \frac{1}{z - M_0} \right\rangle_c \\ &+ \frac{1}{2} \left\langle \text{tr} \frac{V_2'(z_2) - V_2'(M_2)}{z_2 - M_2} \frac{V_1'(z_1) - V_1'(M_1)}{z_1 - M_1} \frac{1}{z_0 - M_0} \text{tr} \frac{1}{z - M_0} \right\rangle_c \\ &- \left\langle \text{tr} \frac{1}{z_0 - M_0} \text{tr} \frac{1}{z - M_0} \right\rangle_c \end{aligned} \quad (231)$$

and the master loop equation is:

$$(V_0'(Z_0) - Z_1)(V_2'(Z_2) - Z_1) - P(Z_0, Z_1, Z_2) = \frac{1}{N^2} U(Z_0, Z_1, Z_2; Z_0) \quad (232)$$

with  $Z_0 = z_0$ ,  $Z_1(z_0) = V_0'(z_0) - W_0(z_0)$  and  $Z_2(z_0) = V_1'(Z_1) - Z_0$ .

### Large $N$ one loop functions

We define:

$$\tilde{Z}_1(z_2) = V_2'(z_2) - W_2(z_2), \quad \tilde{Z}_0(z_2) = V_1'(\tilde{Z}_1(z_2)) - z_2 \quad (233)$$

$$W_{\hat{0},1}(z_0, z_1) = \text{Pol}_{z_0} V_0'(z_0) W_{0,1}(z_0, z_1) \quad (234)$$

$$W_{1,\hat{2}}(z_1, z_2) = \text{Pol}_{z_2} V_2'(z_2) W_{1,2}(z_1, z_2) \quad (235)$$

$$W_{\hat{0},1,\hat{2}}(z_0, z_1, z_2) = \text{Pol}_{z_0, z_2} V_0'(z_0) V_2'(z_2) W(z_0, z_1, z_2) \quad (236)$$

eq.(182) reads:

$$W_{0,2}(z_0, z_2) = \frac{1}{Z_1 - \tilde{Z}_1} \left( \frac{E(z_0, \tilde{Z}_1, z_2)}{z_0 - \tilde{Z}_0} - \frac{E(z_0, Z_1, z_2)}{z_2 - Z_2} \right) \quad (237)$$

eq.(169) reads:

$$U_2(z_0, z_2) = \text{Pol}_{z_2} V'_2(z_2) W_{0,2}(z_0, z_2) = \frac{E(z_0, Z_1, z_2)}{z_2 - Z_2} \quad (238)$$

eq.(172) reads:

$$\begin{aligned} U(z_0, z_1, z_2) &= P_{1,2}(z_0, z_1, z_2) \\ &= V'_2(z_2) - z_1 + \frac{E(z_0, z_1, z_2) - E(z_0, Z_1, z_2)}{z_1 - Z_1} \\ &\quad + \frac{V'_1(z_1) - V'_1(Z_1)}{z_1 - Z_1} \frac{E(z_0, Z_1, z_2)}{z_2 - Z_2} \end{aligned} \quad (239)$$

eq.(32) implies:

$$\begin{aligned} P_{1,1} &= \frac{V'_1(z_1) - V'_1(Z_1)}{z_1 - Z_1} W_{0,2}(z_0, z_2) - \frac{P_{1,2}(z_0, z_1, z_2) - P_{1,2}(z_0, Z_1, z_2)}{z_1 - Z_1} \\ &= \frac{(V'_1 - z_0 - z_2)}{(Z_1 - \tilde{Z}_1)} \left( \frac{E(z_0, \tilde{Z}_1, z_2)}{(z_1 - \tilde{Z}_1)(z_0 - \tilde{Z}_0)} - \frac{E(z_0, Z_1, z_2)}{(z_1 - Z_1)(z_2 - Z_2)} \right) \\ &\quad + 1 - \frac{E(z_0, z_1, z_2)}{(z_1 - Z_1)(z_1 - \tilde{Z}_1)} \end{aligned} \quad (240)$$

We have the relationships:

$$(z_1 - \tilde{Z}_1) P_{0,0}(z_0, z_1, z_2) = \frac{E(z_0, \tilde{Z}_1, z_2)}{z_0 - \tilde{Z}_0} - W_{0,1,\hat{2}}(z_0, z_1, z_2) \quad (241)$$

$$\begin{aligned} (Z_1 - \tilde{Z}_1) W(z_0, z_1, z_2) &= \frac{E(z_0, \tilde{Z}_1, z_2)}{(z_1 - \tilde{Z}_1)(z_0 - \tilde{Z}_0)} - \frac{E(z_0, Z_1, z_2)}{(z_1 - Z_1)(z_2 - Z_2)} \\ &\quad + \frac{(Z_1 - \tilde{Z}_1) W_{0,1,\hat{2}}(z_0, z_1, z_2)}{(z_1 - Z_1)(z_1 - \tilde{Z}_1)} \end{aligned} \quad (242)$$

But so far, we have not been able to compute  $W_{0,1,\hat{2}}$ . We conjecture:

$$\begin{aligned} (z_0 + z_2 - V'_1(z_1)) W_{0,1,\hat{2}}(z_0, z_1, z_2) &= E(z_0, z_1, z_2) \\ - (V'_0(z_0) - z_1 - W_{0,1}(z_0, z_1)) (V'_2(z_2) - z_1 - W_{1,\hat{2}}(z_1, z_2)) \end{aligned} \quad (243)$$

indeed, both sides are polynomials in  $z_0$  and  $z_2$  with the same degree and same large  $z_0$ ,  $z_1$  and  $z_2$  behaviours.

## 9 Conclusion

We have written the loop equations in a very explicit way. To leading order the loop equations become algebraic, and using that fact, we have derived many observables, like the free energy and many correlators.

We have computed some mixed correlators and we have given conjecture for others. So far, the complete one-loop function has not been computed, and we don't even have a conjecture (except in the 3-matrix model). The mixed correlators are important in the study of boundary operators [20].

It would also be interesting to generalize [10] to the chain of matrices, and find the next to leading large  $N$  corrections to the free energy.

*Acknowledgements* The author would like to thank V. Kazakov for encouraging him to translate in english (and add new results to) old notes written in french, and to thank M. Bertola for fruitful discussions and careful reading. When this work was in progress, M. Bertola found independent proofs for some of the formulae in section 5, which will be published later.

## Appendix A Some usefull formulas with residues and polynomial parts

We make an abundant use of the following formulas:

$$1 = \text{Pol}_z \frac{z}{z-Z} = \text{Pol}_z \frac{z}{z-M} = -\text{Res}_z \frac{dz}{z-Z} = -\text{Res}_z \frac{zdz}{(z-Z)(z-M)} \quad (1)$$

$$\text{Pol}_z \frac{V'(z)}{z-Z} = \frac{V'(z) - V'(Z)}{z-Z}, \quad \text{Res}_z \frac{V'(z)dz}{z-Z} = -V'(Z) \quad (2)$$

$$\text{Pol}_z \frac{V'(z)}{(z-Z)(z-M)} = \frac{1}{z-Z} \left( \frac{V'(z) - V'(M)}{z-M} - \frac{V'(Z) - V'(M)}{Z-M} \right) \quad (3)$$

$$\text{Res}_z \frac{V'(z)dz}{(z-Z)(z-M)} = -\frac{V'(Z) - V'(M)}{Z-M} \quad (4)$$

$$\frac{V'(M)}{z-M} = \frac{V'(z)}{z-M} - \frac{V'(z) - V'(M)}{z-M} = \frac{V'(z)}{z-M} - \text{Pol}_z \frac{V'(z)}{z-M} \quad (5)$$

$$\frac{M}{z-M} = \frac{z}{z-M} - 1 = \frac{z}{z-M} - \text{Pol}_z \frac{z}{z-M} \quad (6)$$

## Appendix B Determination of $W_{0,1}$

Define the following polynomial of  $z_0$ :

$$W_{\hat{0},1}(z_0, z_1) = \text{Pol}_{z_0} V'_0(z_0) W_{0,1}(z_0, z_1) \quad (1)$$

eq.(32) implies the following identity (to large  $N$  leading order):

$$(z_1 - Z_1)(1 - W_{0,1}(z_0, z_1)) = z_1 - V'_0(z_0) + W_{\hat{0},1}(z_0, z_1) \quad (2)$$

which can also be written:

$$\left(1 - \frac{Z_1}{z_1}\right)(1 - W_{0,1}(z_0, z_1)) = 1 - \frac{V'_0(z_0) - W_{\hat{0},1}(z_0, z_1)}{z_1} \quad (3)$$

Notice that the RHS is a polynomial in  $z_0$ .

take the log:

$$\ln\left(1 - \frac{Z_1}{z_1}\right) = -\ln(1 - W_{0,1}(z_0, z_1)) + \ln\left(1 - \frac{V'_0(z_0) - W_{\hat{0},1}(z_0, z_1)}{z_1}\right) \quad (4)$$

take the fractionary part ( $\text{Frac } f := f - \text{Pol } f$ ), we find:

$$-\ln(1 - W_{0,1}(z_0, z_1)) = \text{Frac}_{z_0} \left( \ln\left(1 - \frac{Z_1(z_0)}{z_1}\right) \right) \quad (5)$$

i.e.

$$W_{0,1}(z_0, z_1) = 1 - \exp\left(-\text{Frac}_{z_0} \left( \ln\left(1 - \frac{Z_1(z_0)}{z_1}\right) \right)\right) \quad (6)$$

in other words:

$$\begin{aligned} -\ln(1 - W_{0,1}(z_0, z_1)) &= \frac{1}{2i\pi} \oint_{p \in \mathcal{C}_0} \frac{dZ_0(p)}{z_0 - Z_0(p)} \ln\left(1 - \frac{Z_1(p)}{z_1}\right) \\ &= \frac{1}{2i\pi} \oint_{\mathcal{C}_1} \frac{dZ_1(p)}{Z_1(p) - z_1} \ln\left(1 - \frac{Z_0(p)}{z_0}\right) \\ &= \text{Frac}_{z_1} \left( \ln\left(1 - \frac{Z_0(z_1)}{z_0}\right) \right) \\ &= \ln\left(1 - \frac{Z_0(z_1)}{z_0}\right) - \text{Pol}_{z_1} \left( \ln\left(1 - \frac{Z_0(z_1)}{z_0}\right) \right) \end{aligned} \quad (7)$$

where the second equality is obtained by integration by parts.

In other words we have that:

$$(1 - W_{0,1}(z_0, z_1))(z_1 - Z_1(z_0)) = \text{Polynomial in } z_0 \quad (8)$$

$$(1 - W_{0,1}(z_0, z_1)) \prod_{j=1}^{s_1} (z_0 - Z_0(p_{-j,1}(z_1))) = \text{Polynomial in } z_1 \quad (9)$$

that implies that:

$$1 - W_{0,1}(z_0, z_1) = \frac{Q_{0,1}(z_0, z_1)}{(z_1 - Z_1(z_0)) \prod_{j=1}^{s_1} (z_0 - Z_0(p_{-j,1}(z_1)))} \quad (10)$$

where  $Q_{0,1}(z_0, z_1)$  is some polynomial in both variables, of degree  $r_1 + s_1, r_0 + s_0$ , which vanishes each time there exists  $p$  such that  $z_0 = Z_0(p)$  and  $z_1 = Z_1(p)$ . That implies that:

$$\begin{aligned} Q_{0,1}(z_0, z_1) &\propto \prod_{j=1}^{r_1} (z_0 - Z_0(p_{+j,1}(z_1))) \prod_{j=1}^{s_1} (z_0 - Z_0(p_{-j,1}(z_1))) \\ &\propto \prod_{j=1}^{r_0} (z_1 - Z_1(p_{+j,0}(z_0))) \prod_{j=1}^{s_0} (z_1 - Z_1(p_{-j,0}(z_0))) \end{aligned} \quad (11)$$

i.e.  $Q_{0,1}$  must be equal to (same degree, same zeroes):

$$Q_{0,1}(z_0, z_1) \propto \oint dz_2 \dots dz_N \frac{E(z_0, z_1, z_2, \dots, z_N)}{\prod_{j=1}^{N-1} (z_{j+1} + z_{j-1} - V_j(z_j))} \quad (12)$$

We thus have:

$$W_{0,1}(z_0, z_1) = 1 - \frac{Q_{0,1}(z_0, z_1)}{(z_1 - Z_1(z_0)) \prod_{j=1}^{s_1} (z_0 - Z_0(p_{-j,1}(z_1)))}$$

(13)

i.e.

$$W_{0,1}(z_0, z_1) = 1 + g_{0,d_0+1} \frac{\prod_{j=1}^{r_1} (z_0 - Z_0(p_{+j,1}(z_1)))}{(z_1 - Z_1(p_{+1,0}(z_0)))}$$

(14)

## References

- [1] J. Ambjorn, L. Chekhov, C.F. Kristjansen, Yu. Makeenko, “Matrix Model Calculations beyond the Spherical Limit”, *Nucl.Phys.* **B404** (1993) 127-172. Erratum-ibid. **B449** (1995) 681.
- [2] M. Bertola, “Bilinear semi-classical moment functionals and their integral representation”, *J. App. Th.* **121** 71–99 (2003).
- [3] M. Bertola, ”Free Energy of the Two-Matrix Model/dToda Tau-Function”, preprint CRM-2921 (2003), hep-th/0306184.
- [4] M. Bertola, B. Eynard, J. Harnad, “Differential systems for biorthogonal polynomials appearing in 2-matrix models and the associated Riemann-Hilbert problem”, preprint CRM-2852 (2002), Saclay T02/097, nlin.SI/0208002.
- [5] Duality, M. Bertola, B.Eynard, J. Harnad “Bi-orthogonal Polynomials and Multi-Matrix Models”, *Commun.Math.Phys.* **229** (2002) 73-120, xxx, nlin.SI/0108049.
- [6] G. Bonnet, F. David, B. Eynard, “Breakdown of universality in multi-cut matrix models”, *J.Phys. A* **33** 6739-6768 (2000).
- [7] J.M. Daul, V. Kazakov, I.K. Kostov, “Rational Theories of 2D Gravity from the Two-Matrix Model”, *Nucl. Phys.* **B409**, 311-338 (1993), hep-th/9303093.
- [8] P. Di Francesco, P. Ginsparg, J. Zinn-Justin, “2D Gravity and Random Matrices”, *Phys. Rep.* **254**, 1 (1995).
- [9] N. M. Ercolani and K. T.-R. McLaughlin “Asymptotics and integrable structures for biorthogonal polynomials associated to a random two-matrix model”, *Physica D*, *152-153*, 232-268 (2001).
- [10] B. Eynard, “Large N expansion of the 2-matrix model”, *JHEP* **01** (2003) 051, hep-th/0210047.
- [11] B. Eynard, “Large N expansion of the 2-matrix model, multicut case”, preprint SPHT03/106, ccsc-00000521, math-ph/0307052.

- [12] B. Eynard, “Eigenvalue distribution of large random matrices, from one matrix to several coupled matrices” *Nucl. Phys. B***506**, 633 (1997), cond-mat/9707005.
- [13] B. Eynard, “Correlation functions of eigenvalues of multi-matrix models, and the limit of a time dependent matrix”, *J. Phys. A: Math. Gen.* **31**, 8081 (1998), cond-mat/9801075.
- [14] B. Eynard “An introduction to random matrices”, lectures given at Saclay, October 2000, notes available at <http://www-spht.cea.fr/articles/t01/014/>.
- [15] B. Eynard “Gravitation quantique bidimensionnelle et matrices aléatoires”, Thèse de doctorat de l’université Paris 6 (1995).
- [16] H.M. Farkas, I. Kra, ”Riemann surfaces” 2nd edition, Springer Verlag, 1992.
- [17] J.D. Fay, ”Theta functions on Riemann surfaces”, Springer Verlag, 1973.
- [18] A. Guionnet, Zeitouni, “Large deviations asymptotics for spherical integrals”, *J. F. A.* **188**, 461–515 (2002).
- [19] V.A. Kazakov, A. Marshakov, ”Complex Curve of the Two Matrix Model and its Tau-function”, *J.Phys.* **A36** (2003) 3107-3136, hep-th/0211236.
- [20] I. K. Kostov, “Boundary Correlators in 2D Quantum Gravity: Liouville versus Discrete Approach”, *Nucl.Phys.* **B658** (2003) 397-416, hep-th/0212194.
- [21] M.L. Mehta, *Random Matrices*, 2nd edition, (Academic Press, New York, 1991).
- [22] M. Staudacher, “Combinatorial solution of the 2-matrix model”, *Phys. Lett.* **B305** (1993) 332-338.